

# Performance of Lagrangian Finite Volume Approaches for Linear and Nonlinear Solid Mechanics Analyses

Philip Cardiff

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Vila Flor Cultural Centre  
Guimarães, Portugal

# Performance of Lagrangian Finite Volume Approaches for Linear and Nonlinear Solid Mechanics Analyses

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# Background & Motivation



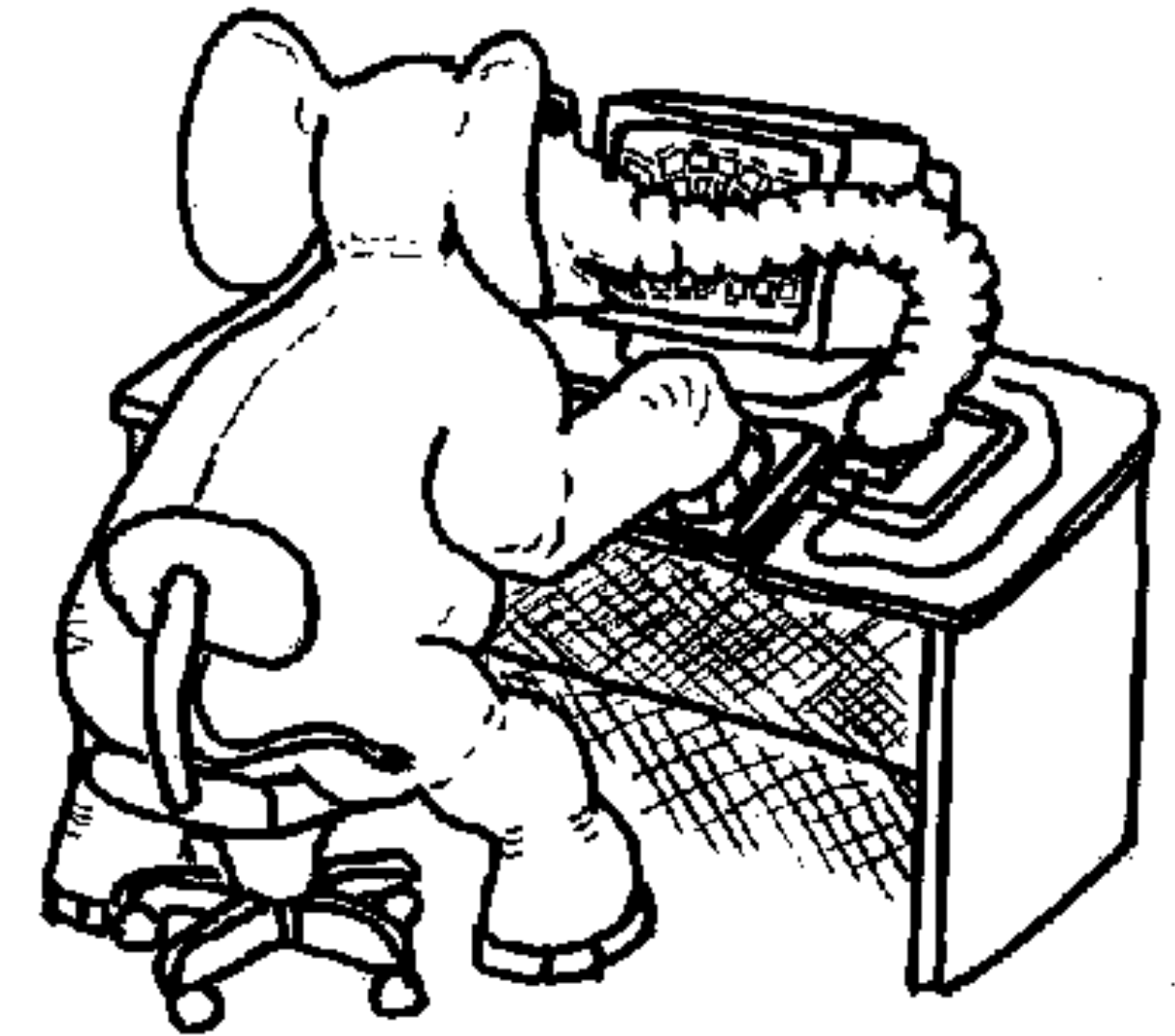
# Background

Historically

**Computational  
Solid Mechanics** = **Finite Element  
Method**

... but alternative approaches, such as the **Finite Volume Method**, are gaining traction ...

F.E.M.



Finite Elephant Modeler

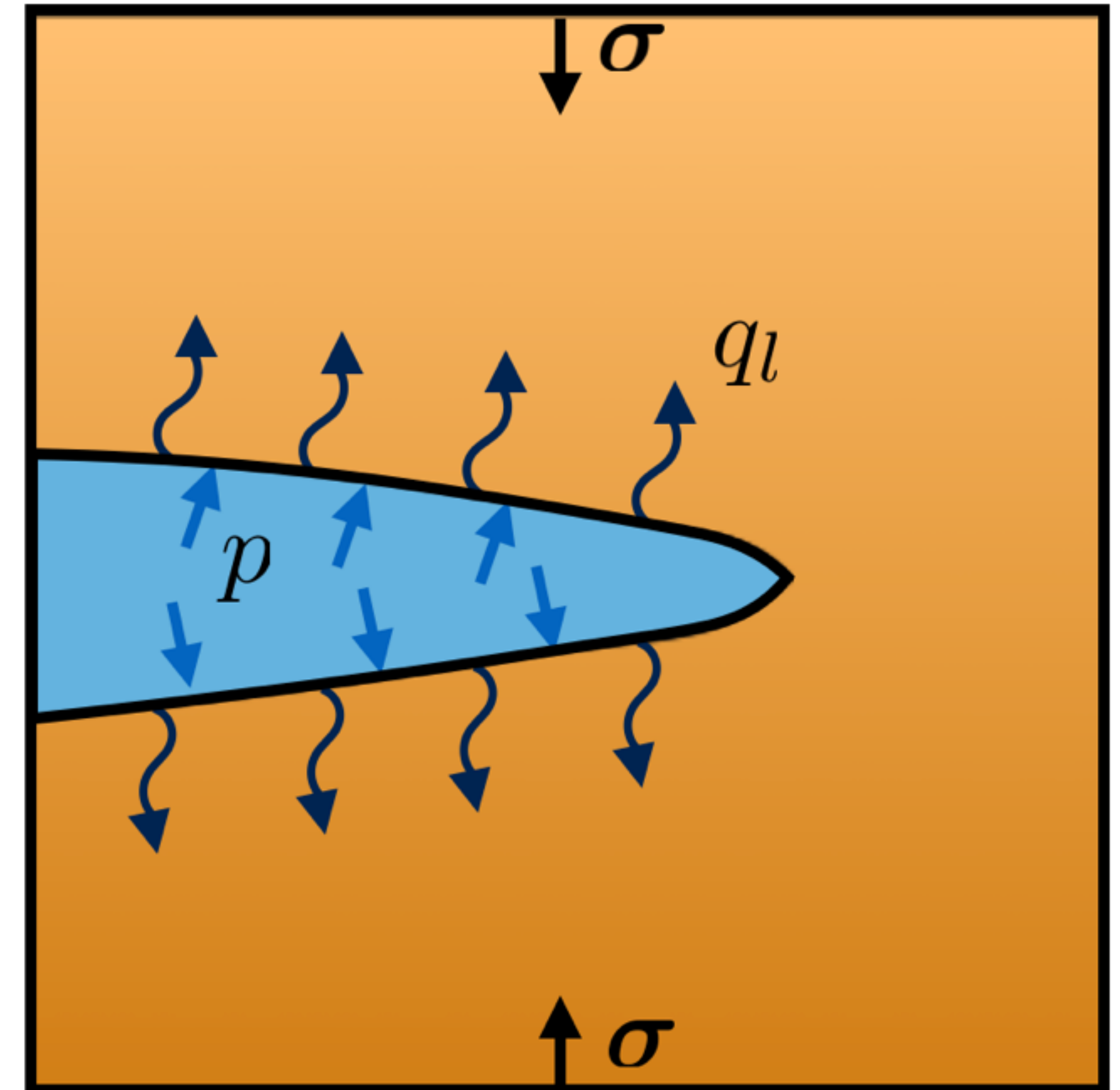
*J. D. Achenbach*  
'92

# Motivation

There are a number of motivations for pursuing FV-based solid mechanics:

- complex nonlinear **multi-physics** applications
- **Fluid-Solid Interaction** within the same numerical framework
- the relative success of OpenFOAM for CFD has led to a desire for equivalent solid mechanics procedures within the library
- **academic interest** ...

**Let's look at some of the latest FV solid mechanics in OpenFOAM.**



$$\oint_{\Gamma} \mathbf{n} \cdot \mathbf{D} \, d\Gamma - \int_{\Omega} Q \, d\Omega = 0$$

$$\rho \frac{\partial^2 u}{\partial t^2} = \nabla \cdot \sigma$$

$$\rho C \frac{\partial T}{\partial t} = k \nabla^2 T$$

# Finite Volume Method for Solid Mechanics

# Mathematical Model

## Conservation of Linear Momentum

$$\underbrace{\frac{D}{Dt} \int_{\Omega} \rho \mathbf{v} \, d\Omega}_{\text{change in momentum}} = \underbrace{\int_{\Gamma} \mathbf{n} \cdot \boldsymbol{\sigma} \, d\Gamma}_{\text{sum of surface forces}} + \underbrace{\int_{\Omega} \rho \mathbf{f}_b \, d\Omega}_{\text{sum of body forces}}$$

Let us consider two contrasting **constitutive relations**:

Hooke's law (small strains/rotations, linear):

$$\boldsymbol{\sigma} = 2\mu\boldsymbol{\epsilon} + \lambda \operatorname{tr}(\boldsymbol{\epsilon})\mathbf{I}$$

Neo-Hookean hyperelastoplastic law (large strains/rotations nonlinear):

$$\boldsymbol{\sigma} = J \left[ \mu \operatorname{dev}(\bar{\mathbf{b}}^e) - \frac{\kappa}{2} (J^2 - 1) \mathbf{I} \right]$$

Calculated from Mises-Huber-Levy  $J_2$  plasticity

Inserting the **constitutive relation** into the **governing equation** gives the mathematical model:

Hookean linear elastic solid:

$$\rho \frac{\partial^2 \mathbf{u}}{\partial t^2} = \nabla \cdot (\mu \nabla \mathbf{u}) + \nabla \cdot (\mu (\nabla \mathbf{u})^T) + \nabla \cdot (\lambda \text{tr}(\nabla \mathbf{u}) \mathbf{I}) + \rho \mathbf{f}_b$$

Neo-Hookean nonlinear hyperelastoplastic solid (updated Lagrangian formulation):

$$\rho \frac{\partial^2 \mathbf{u}}{\partial t^2} = \nabla \cdot \left( j \boldsymbol{\sigma} \cdot \mathbf{f}^{-T} \right) + \rho \mathbf{f}_b$$

## Solution algorithm: block-coupled

Hookean linear elastic solid:

$$\rho \frac{\partial^2 \mathbf{u}}{\partial t^2} = \nabla \cdot (\mu \nabla \mathbf{u}) + \nabla \cdot (\mu (\nabla \mathbf{u})^T) + \nabla \cdot (\lambda \text{tr}(\nabla \mathbf{u}) \mathbf{I}) + \rho \mathbf{f}_b$$

`fvm::laplacian`

`fvm::laplacianTranspose`

`fvm::laplacianTrace`

**... not implemented in OpenFOAM ...**

**... we need to do it ourselves.**

## Solution algorithm: segregated

Neo-Hookean nonlinear hyperelastoplastic solid:

$$\rho \frac{\partial^2 \mathbf{u}}{\partial t^2} = \overbrace{K_{\text{imp}} \nabla^2 (\Delta \mathbf{u})}^{\text{implicit}} + \overbrace{\nabla \cdot (j \boldsymbol{\sigma} \cdot \mathbf{f}^{-T}) - K_{\text{imp}} \nabla^2 (\Delta \mathbf{u})}^{\text{explicit}}$$

$\downarrow$  `fvm::laplacian`                       $\downarrow$  `fvc::div`

# Block-Coupled Challenges:

## Equation

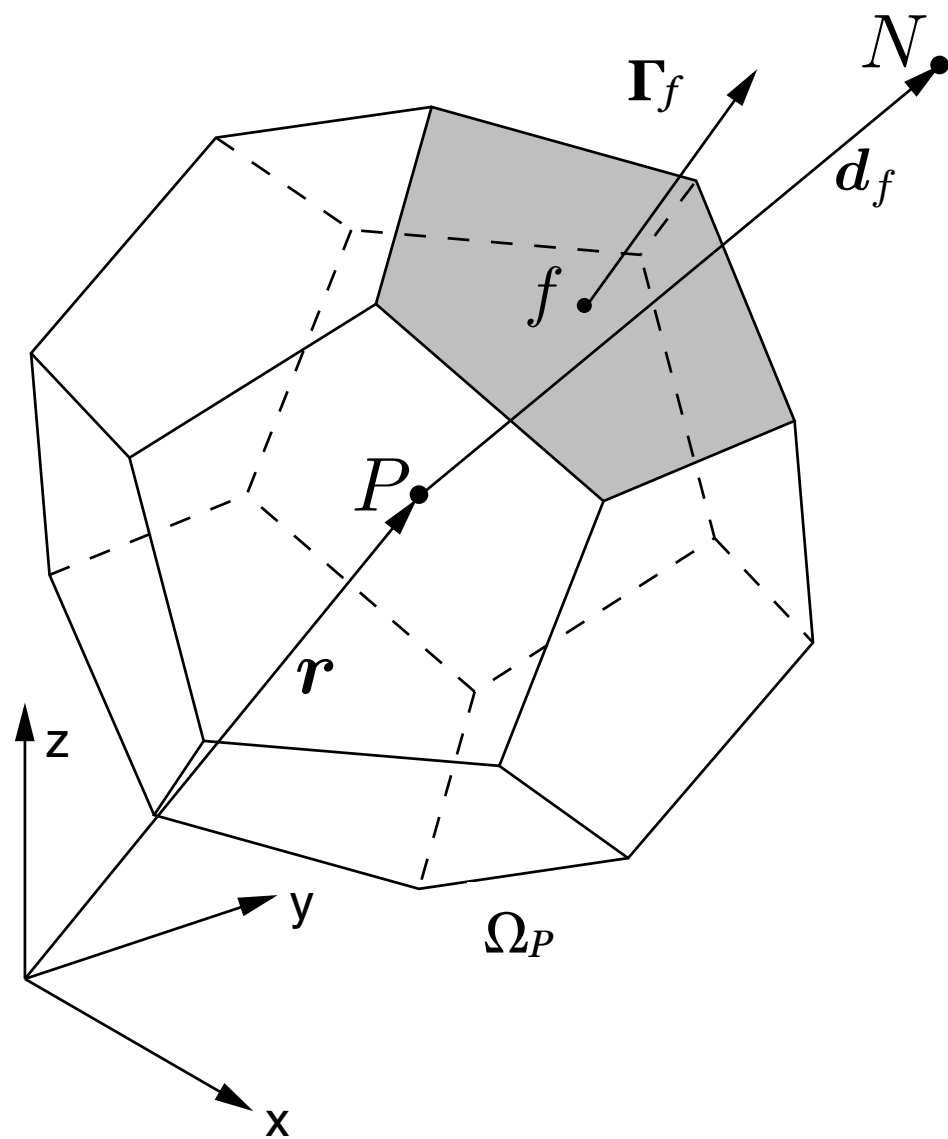
## Discretisation

# Normal Gradient Discretisation

Normal derivatives approximated using central differencing

$$\int_{\Gamma} (2\mu + \lambda) \nabla_n u_n \, d|\Gamma| \approx \sum_f (2\mu_f + \lambda_f) |\Delta_f| (\mathbf{n}_f \mathbf{n}_f) \cdot \left( \frac{\mathbf{u}^N - \mathbf{u}^P}{|d_f|} \right) |\Gamma_f|$$

$$+ \sum_f (2\mu_f + \lambda_f) (\mathbf{n}_f \mathbf{n}_f) \cdot [\mathbf{k}_f \cdot (\nabla_t \mathbf{u})_f] |\Gamma_f|$$

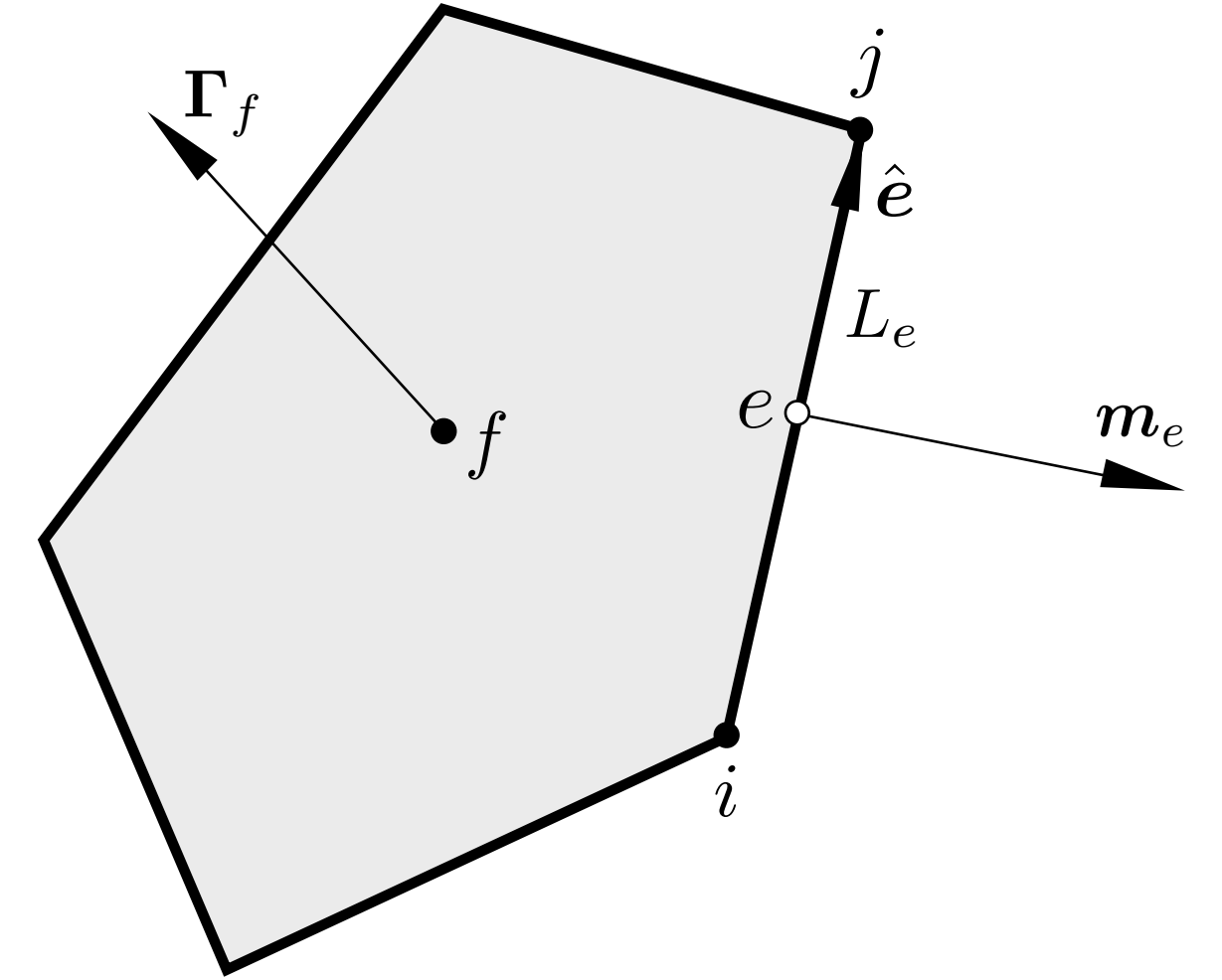


non-orthogonal correction implicitly discretised (see next slide); even gradient term is treated implicitly!

# Tangential Gradient Discretisation

Tangential derivatives approximated using Finite Area Method

$$(\nabla_t u_n)_f = \frac{1}{|\Gamma_f|} \sum_e \mathbf{m}_e (u_n)_e L_e$$



$$\oint_{\Gamma} \mu \nabla_t u_n \, d\Gamma \approx \sum_f \sum_e \sum_{pc} \boxed{\frac{1}{2} \mu_f \omega_{pc} L_e (\mathbf{m}_e \mathbf{n}_f)} \cdot \mathbf{u}^{pc}$$

our coefficients!

$$\oint_{\Gamma} \lambda \mathbf{n} \operatorname{tr}(\nabla_t \mathbf{u}_t) \, d\Gamma \approx \sum_f \sum_e \sum_{pc} \boxed{\frac{1}{2} \lambda_f \omega_{pc} L_e (\mathbf{m}_e \mathbf{n}_f)^T} \cdot \mathbf{u}^{pc}$$

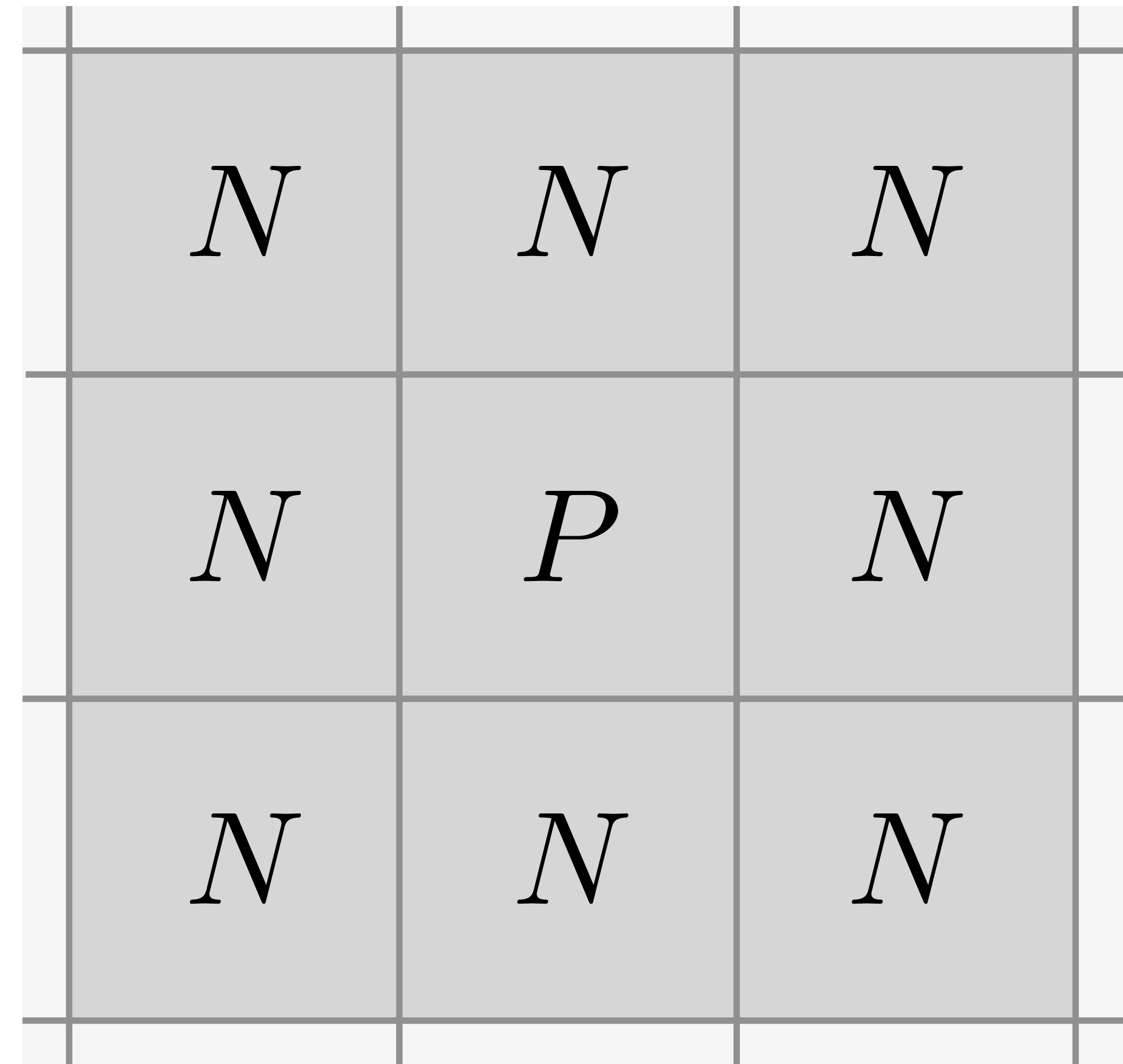
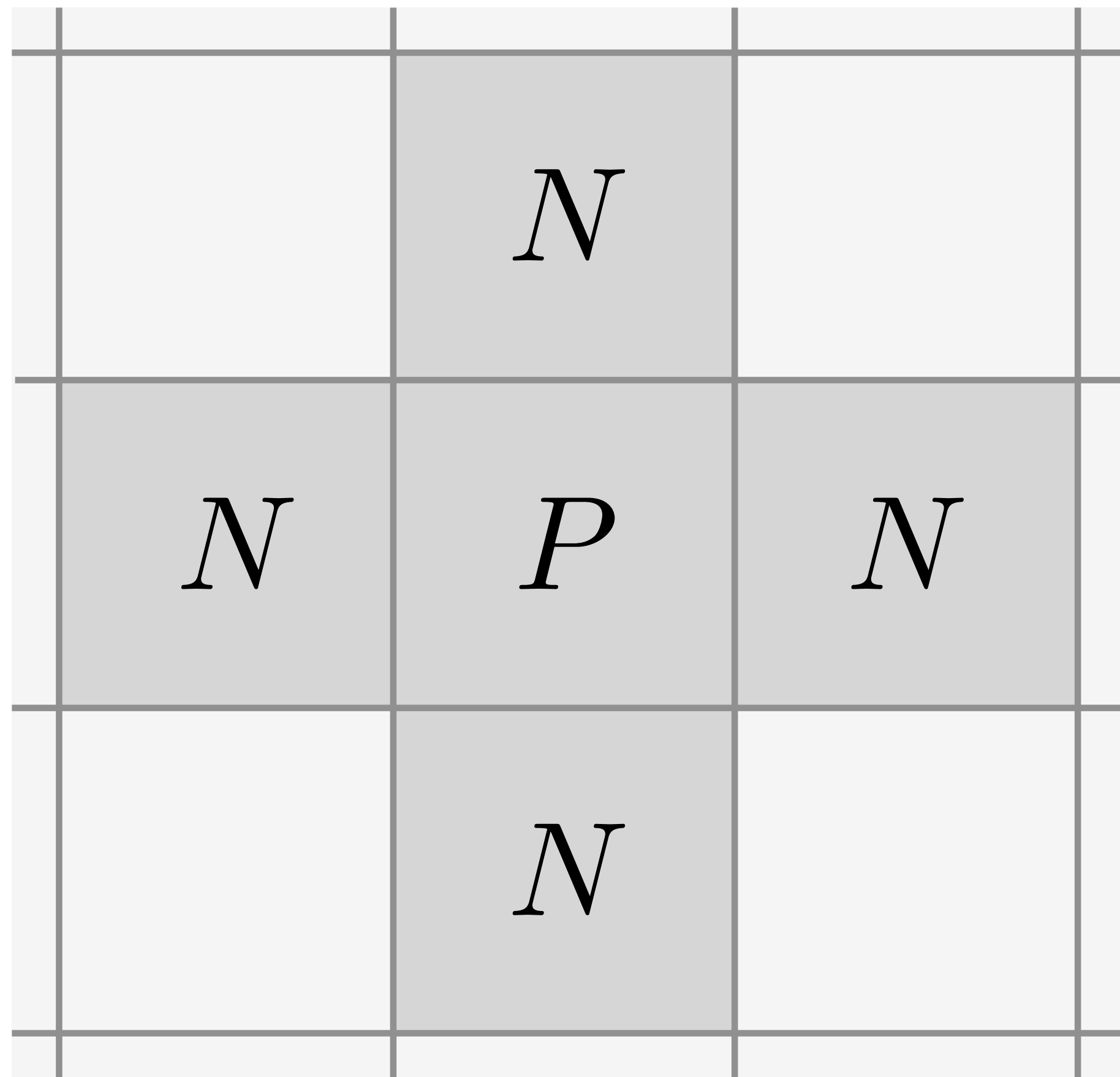
```
solve  
(  
    magic::div(sigma)  
)
```

# Block-Coupled Challenges: Implementation

# Sparse Matrices in OpenFOAM

OpenFOAM fvMesh uses a **sparse matrix addressing** based on the mesh internal faces: only **near-neighbours** may be included implicitly in the matrix.

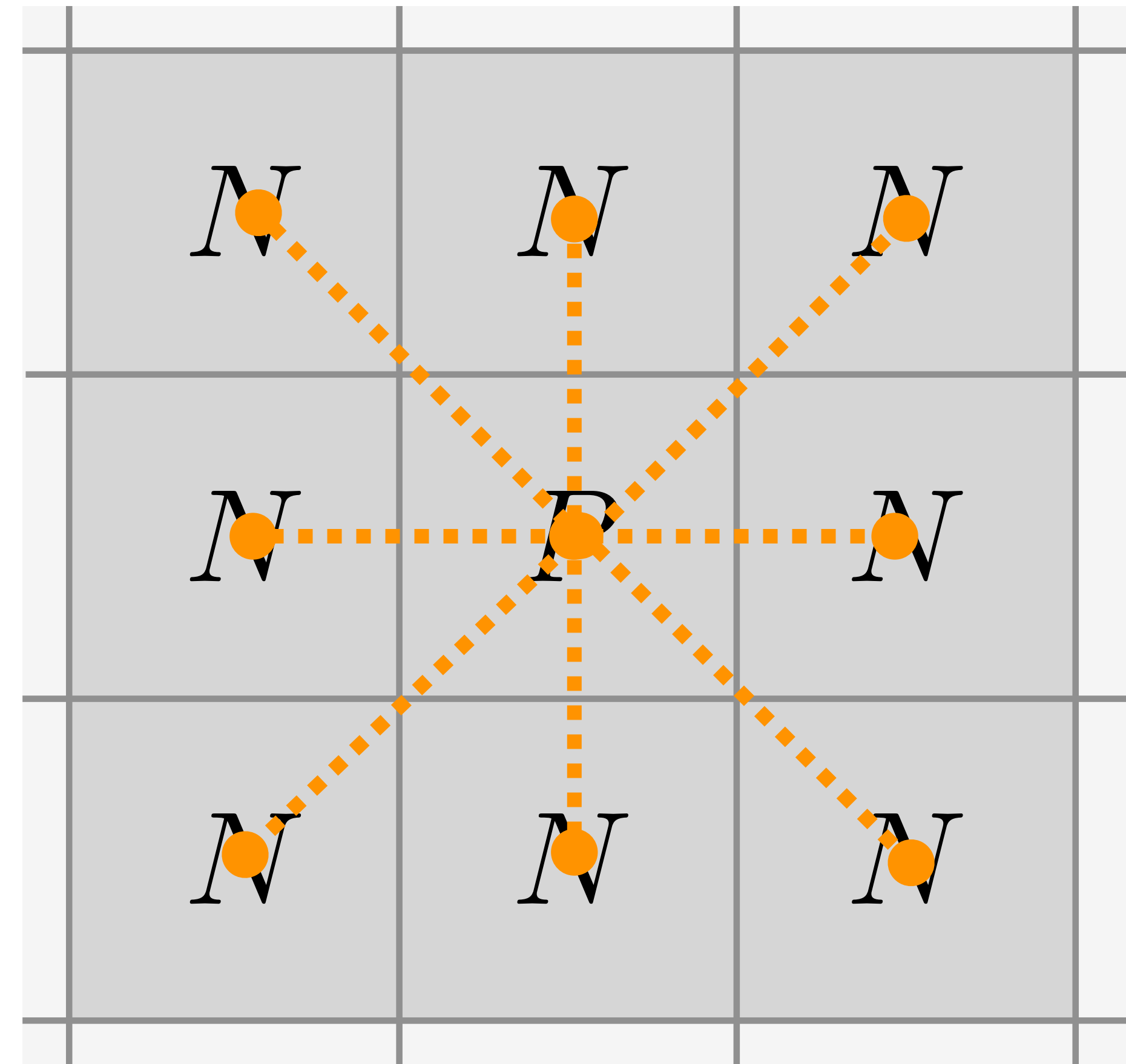
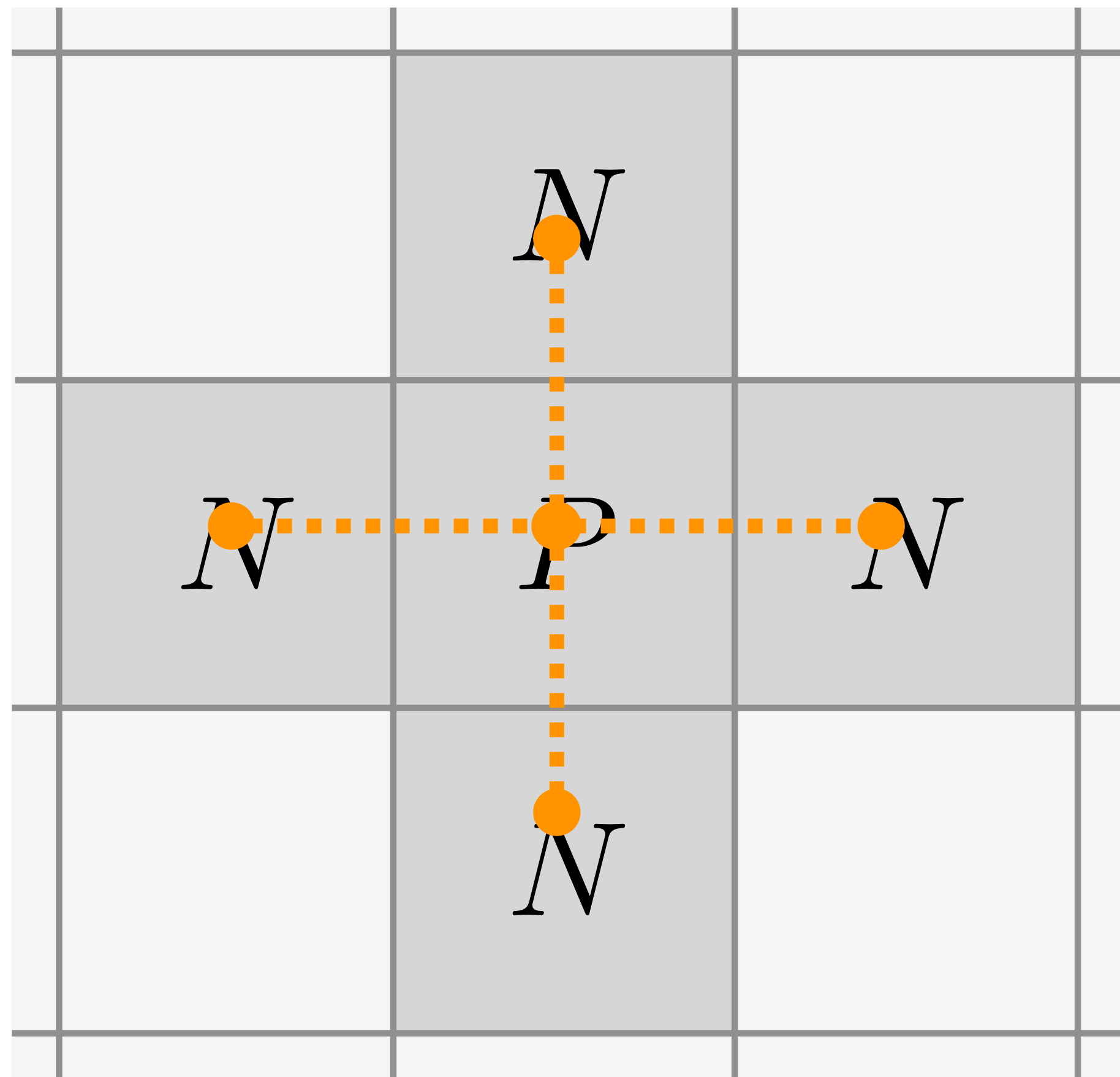
**BUT, we require cell-point-cell neighbours to be included implicitly.**

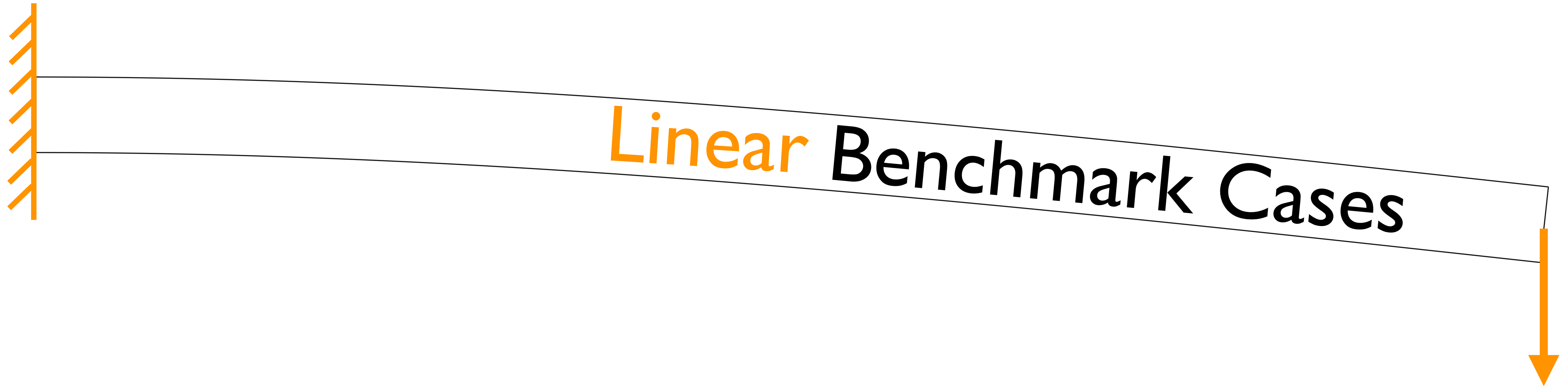


# New Sparse Matrix Addressing Required

Let us refer to the matrix connections as **implicit bonds**

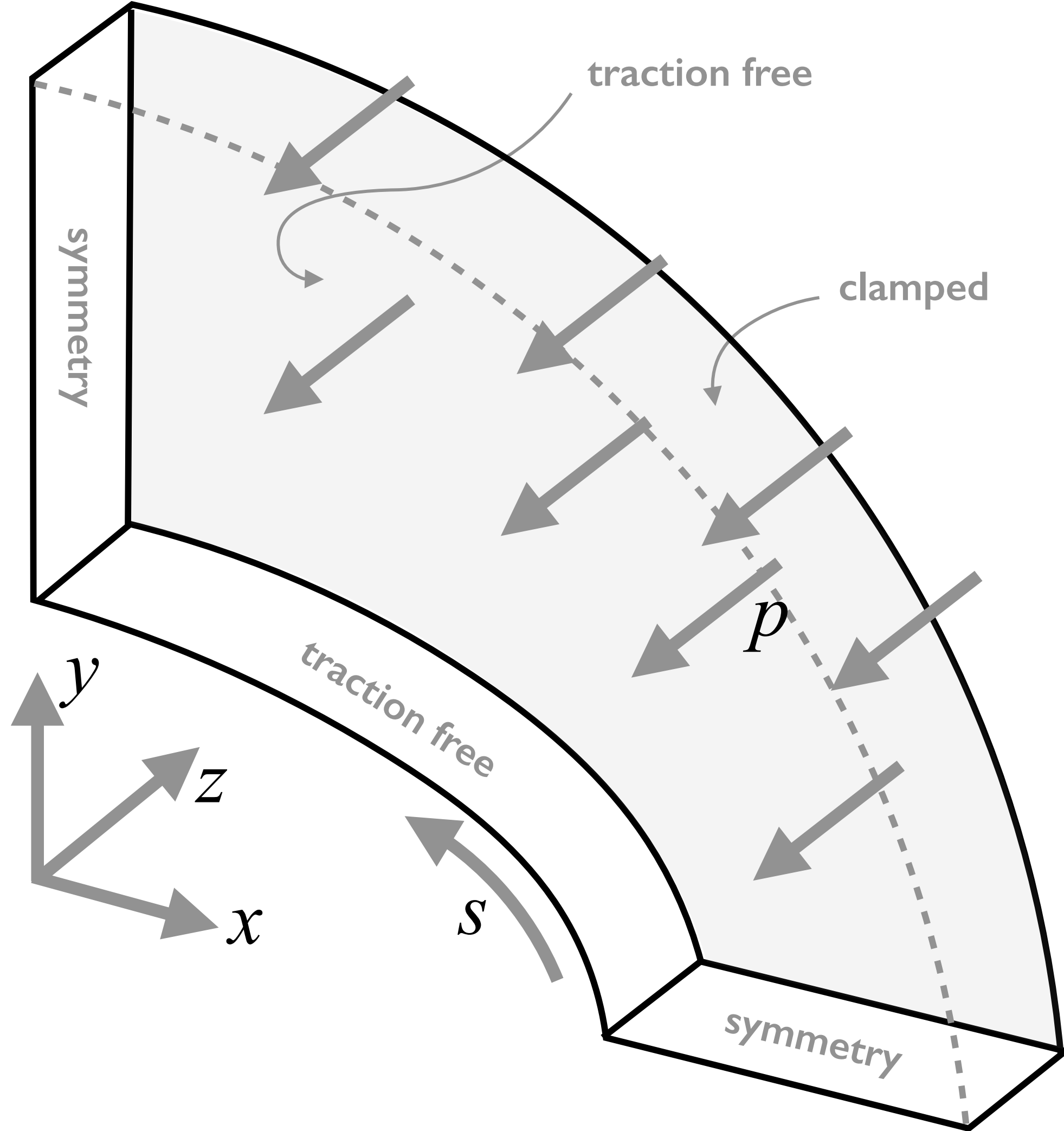
- Standard addressing allows implicit bonds with face neighbours
- New addressing must allow implicit bonds with point neighbours





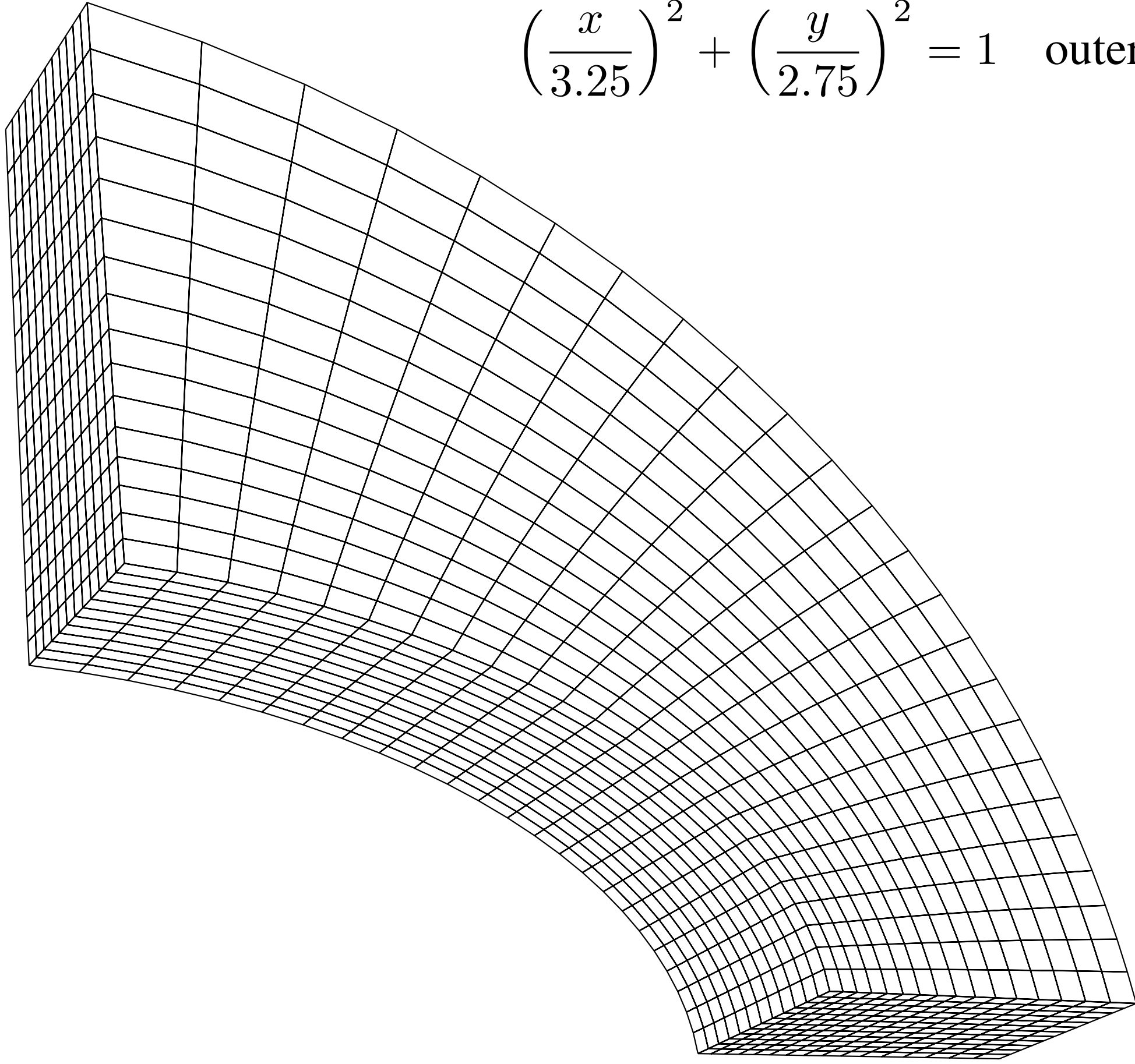
**Linear** Benchmark Cases

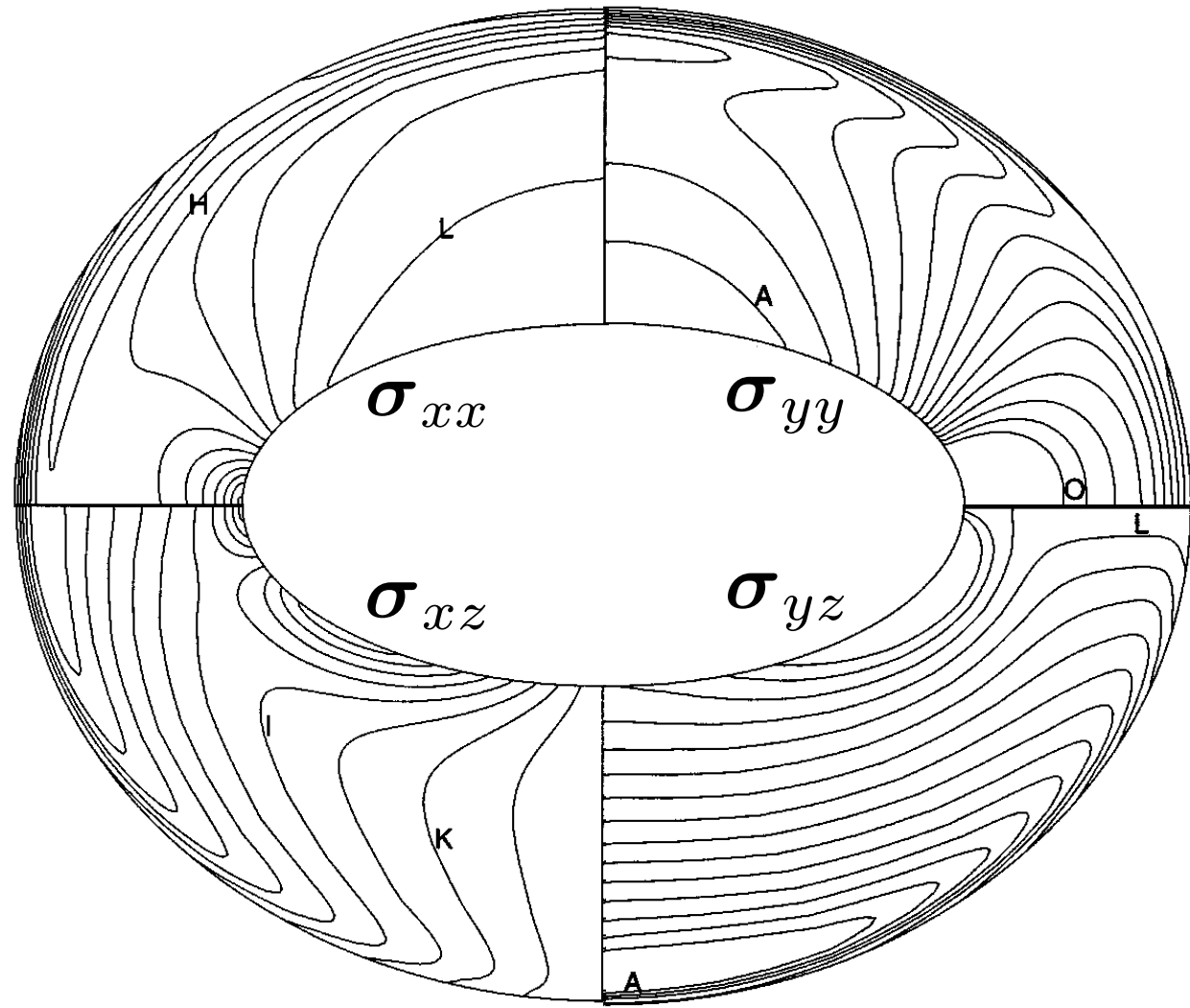
# Out-of-Plane Bending of an Elliptic Plate



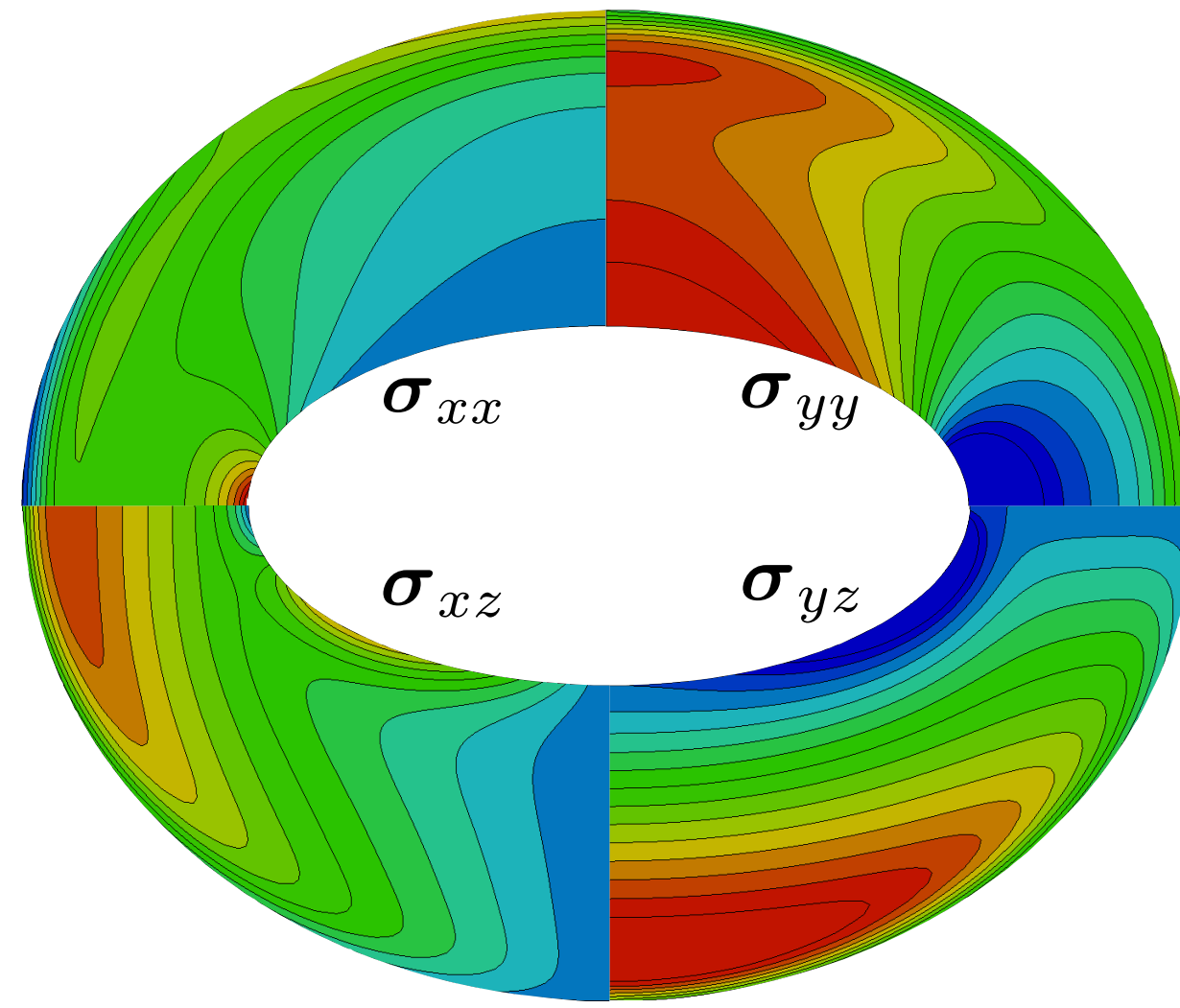
$$\left(\frac{x}{2}\right)^2 + \left(\frac{y}{1}\right)^2 = 1 \quad \text{inner ellipse}$$

$$\left(\frac{x}{3.25}\right)^2 + \left(\frac{y}{2.75}\right)^2 = 1 \quad \text{outer ellipse}$$

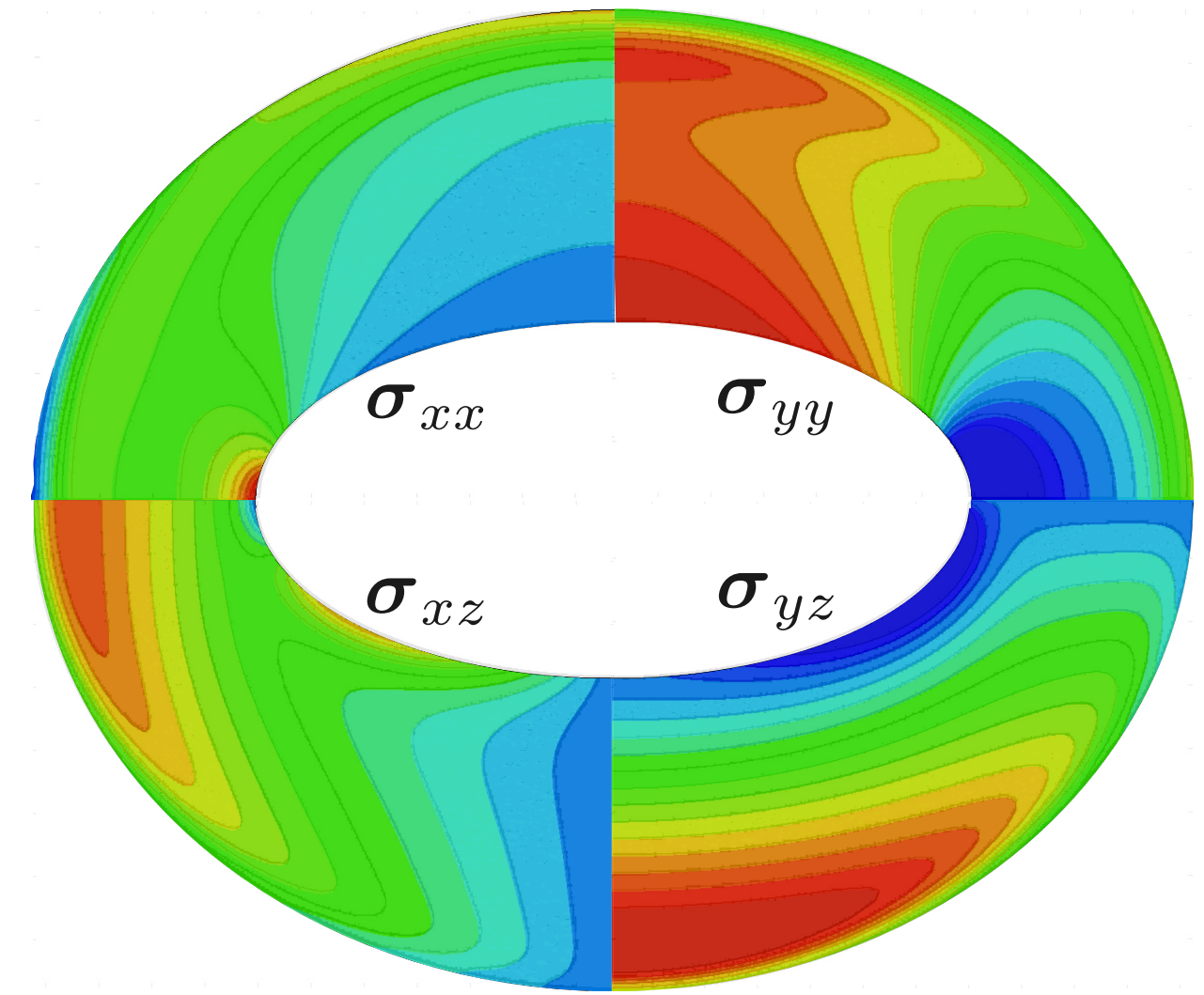




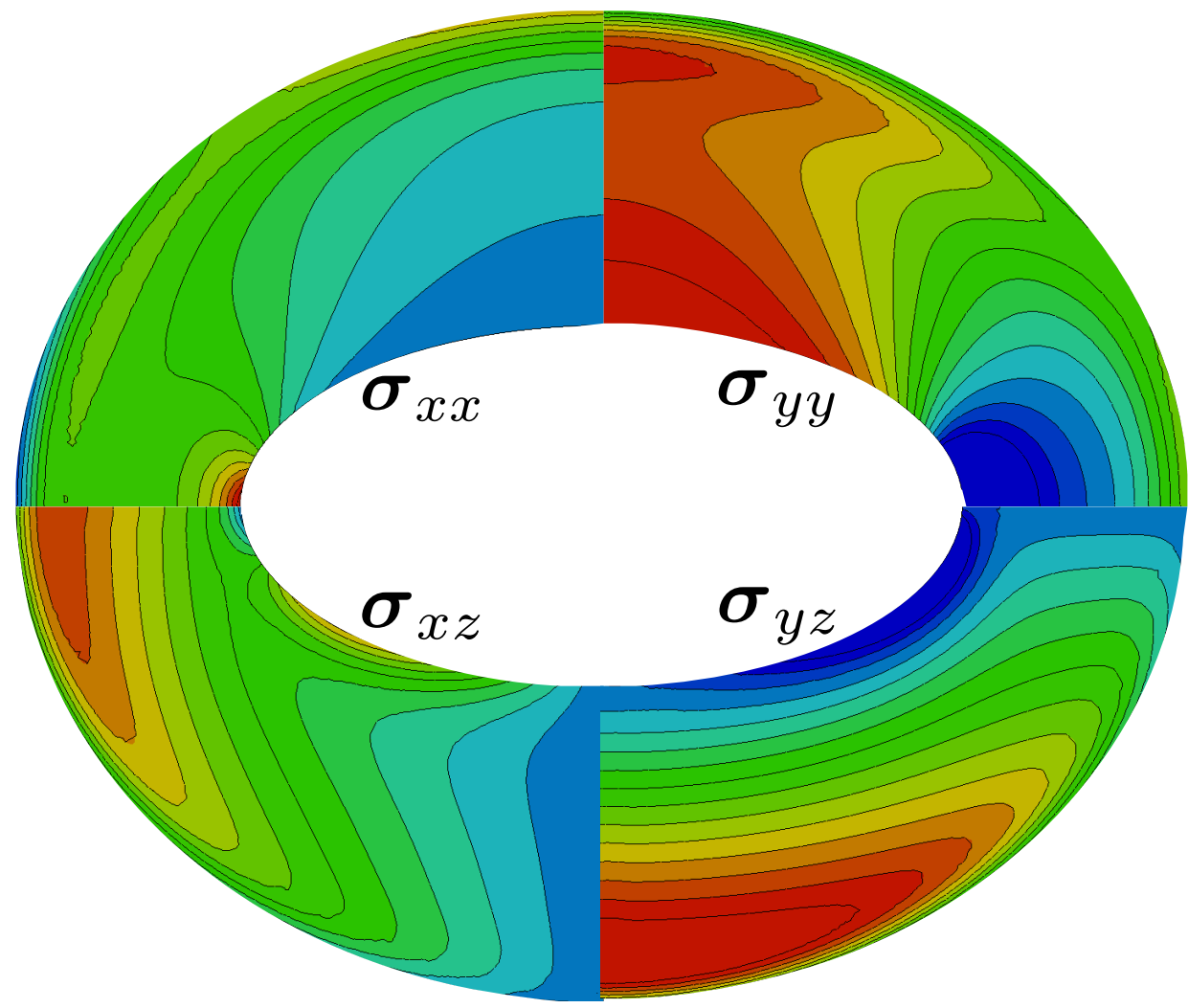
(a) Demirdžić et al. [3] Reference Solution



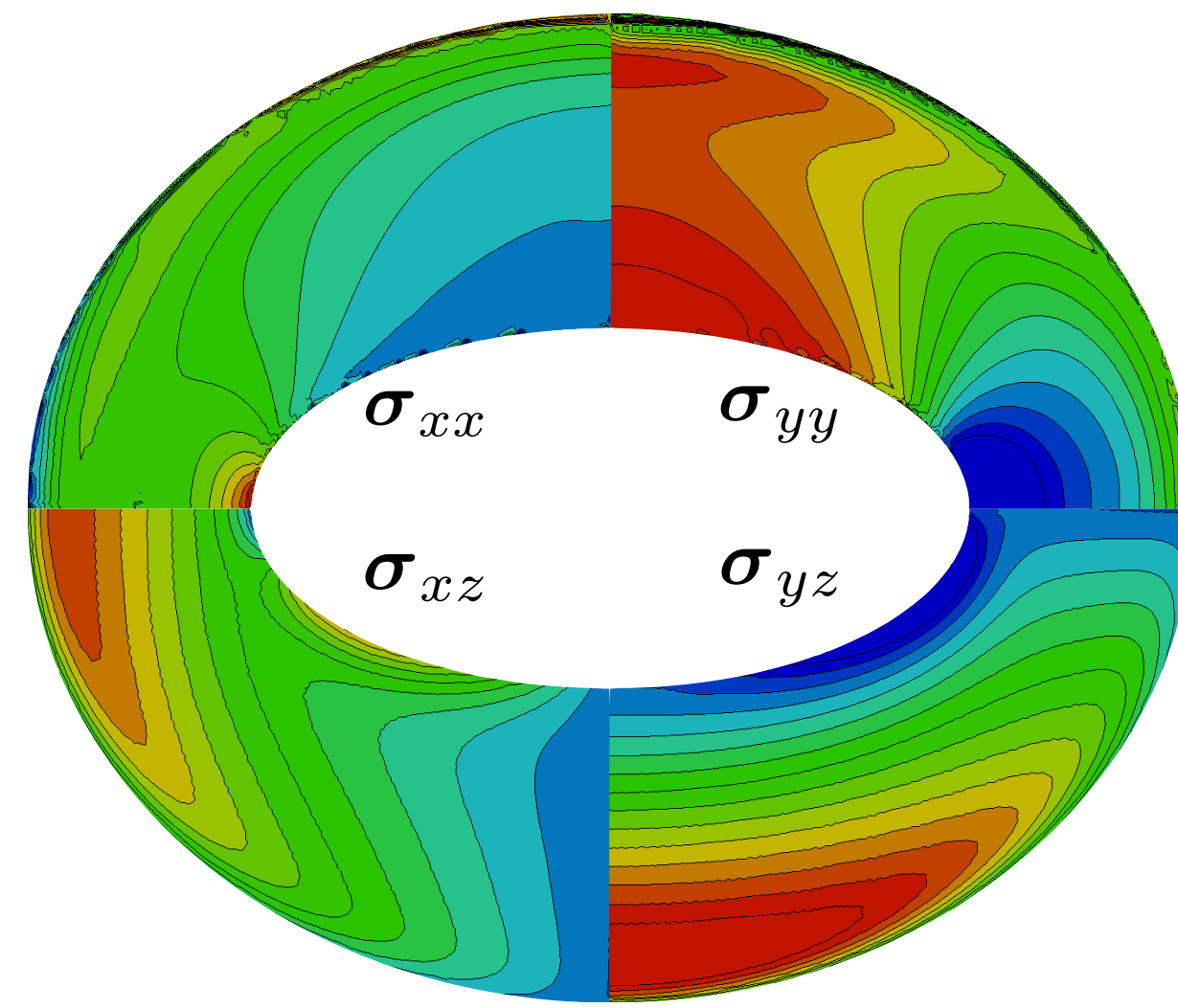
(b) Hexahedral Mesh



(e) Abaqus (Hexahedral Mesh)



(c) Polyhedral Mesh



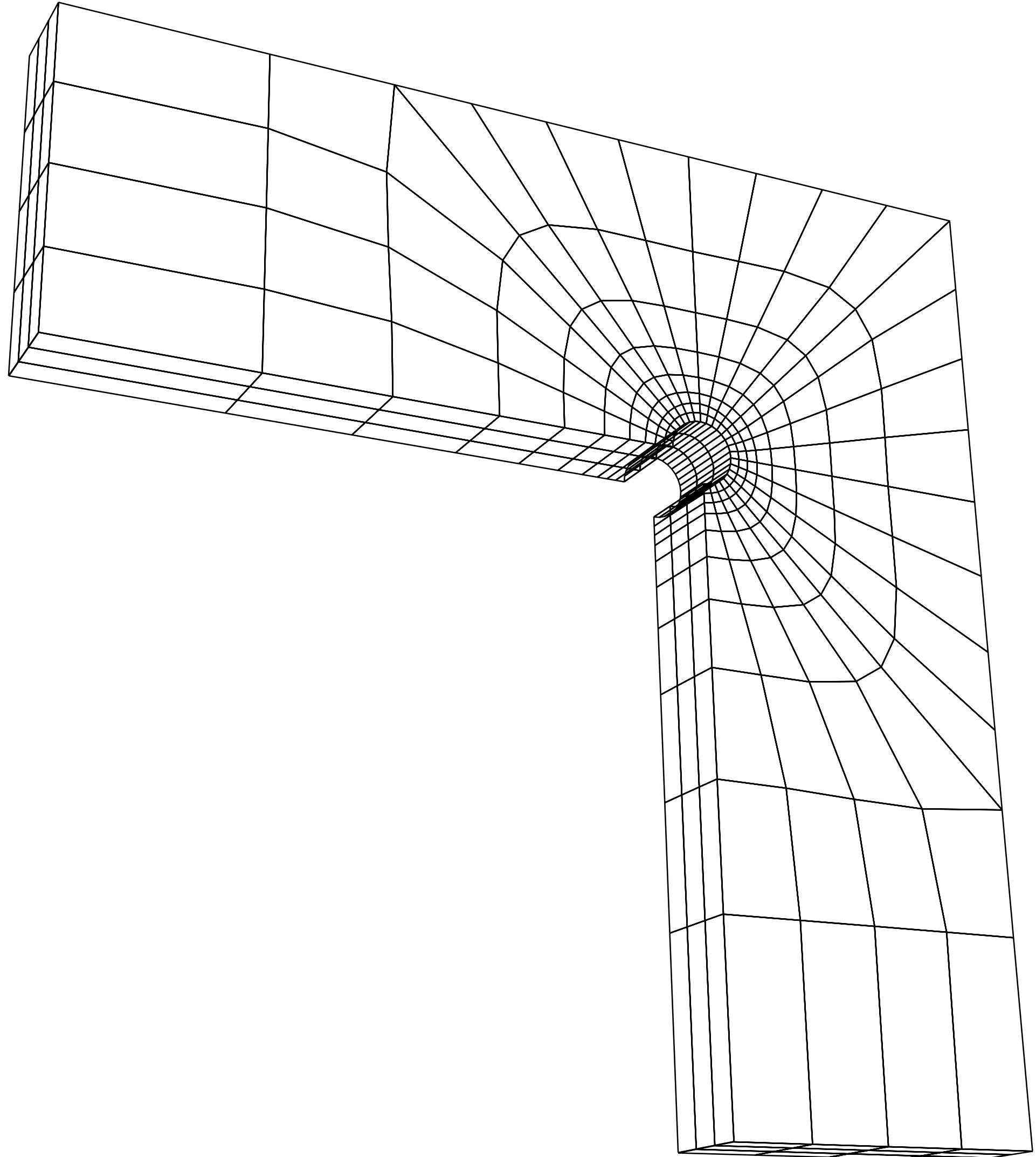
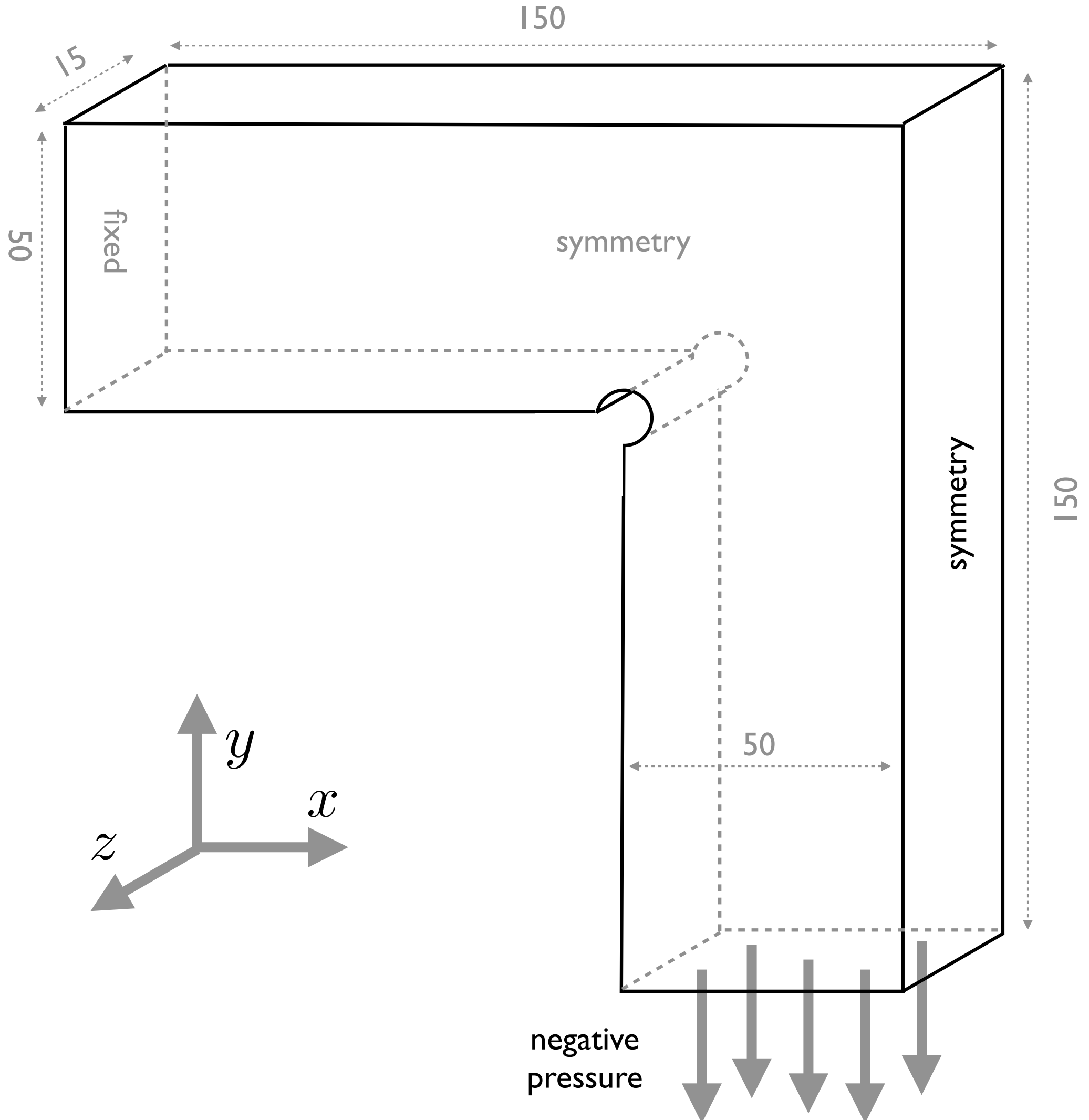
(d) Tetrahedral Mesh

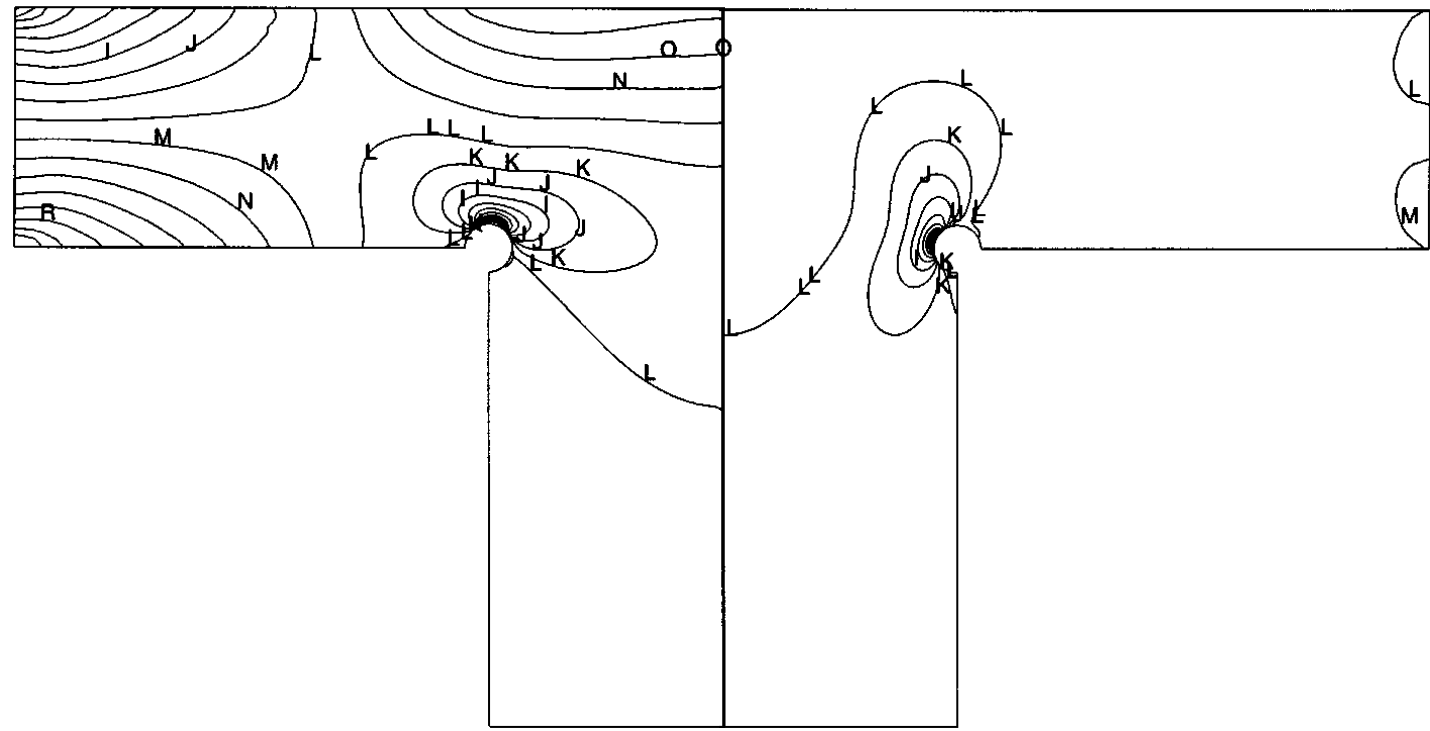
# Case 2: Out-of-Plane Bending of an Elliptic Plate

Mesh	Coupled		Segregated <sup>1</sup>		Segregated <sup>2</sup>		Abaqus	
	Time	Memory	Time	Memory	Time	Memory	Time	Memory
72	0.03	7	0.5	6	0.5	6	4	24
576	0.15	13	1	8	2	9	5	31
4 608	1.6	51	6.5	20	14	30	6	107
36 864	11	300	102	80	145	120	34	1 197
294 912	242	2 200	1 474	500	1 990	800	1 375	17 900

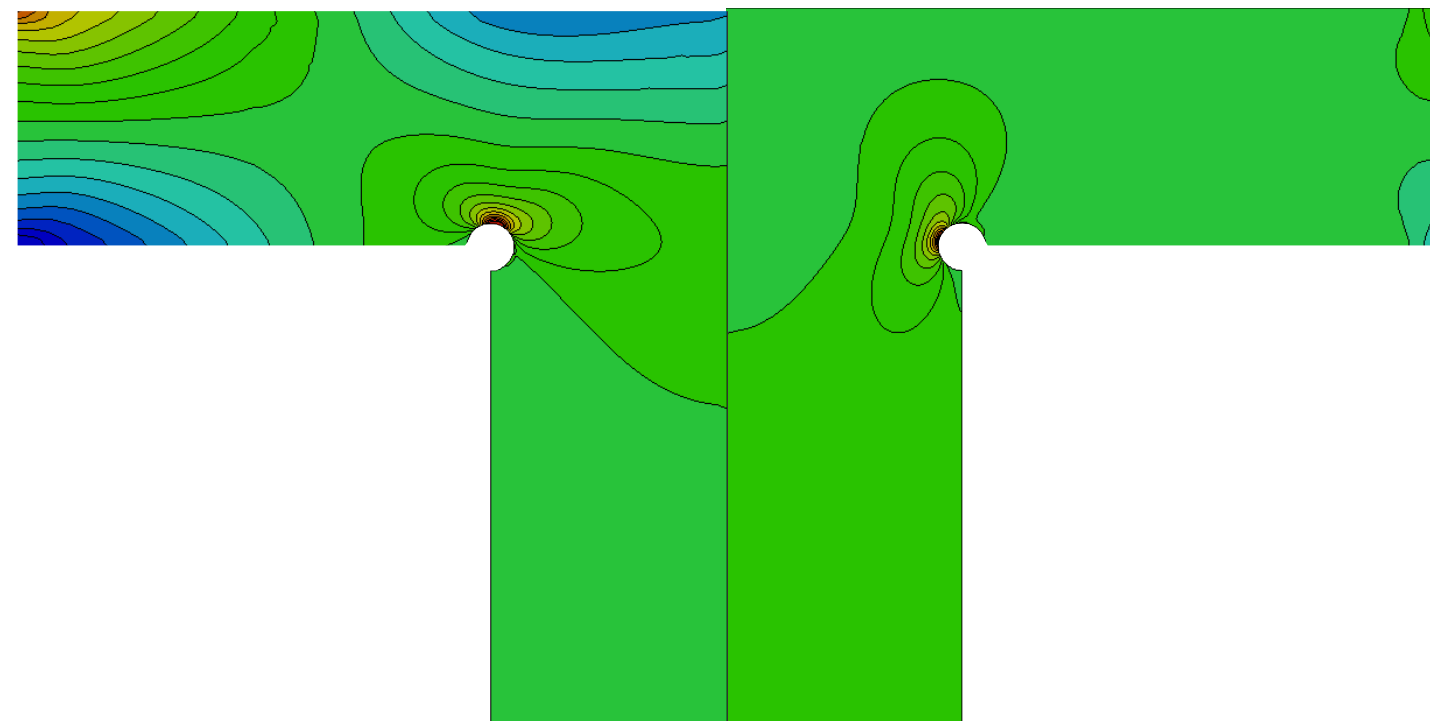
Table II. Elliptic plate: wall-clock time (in s) and maximum memory usage (in MB)

# Narrow T-section Component

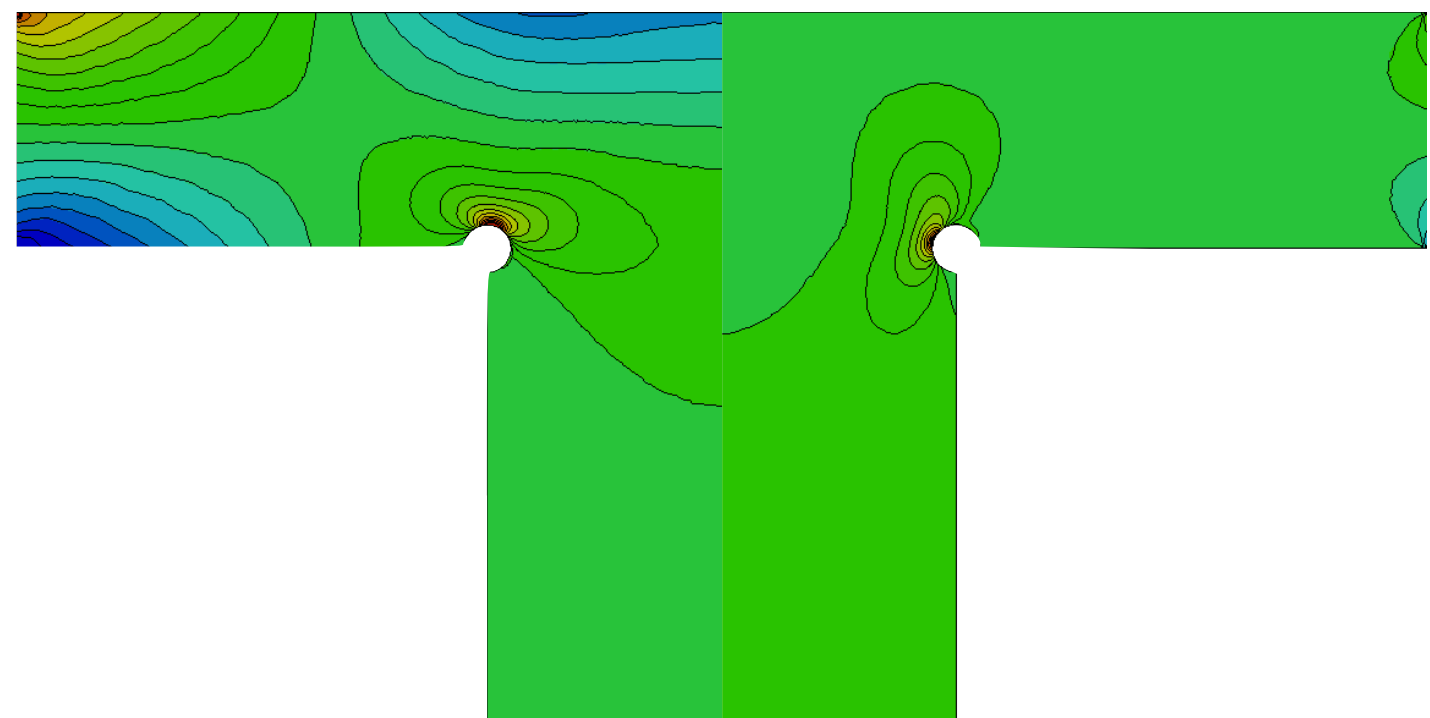




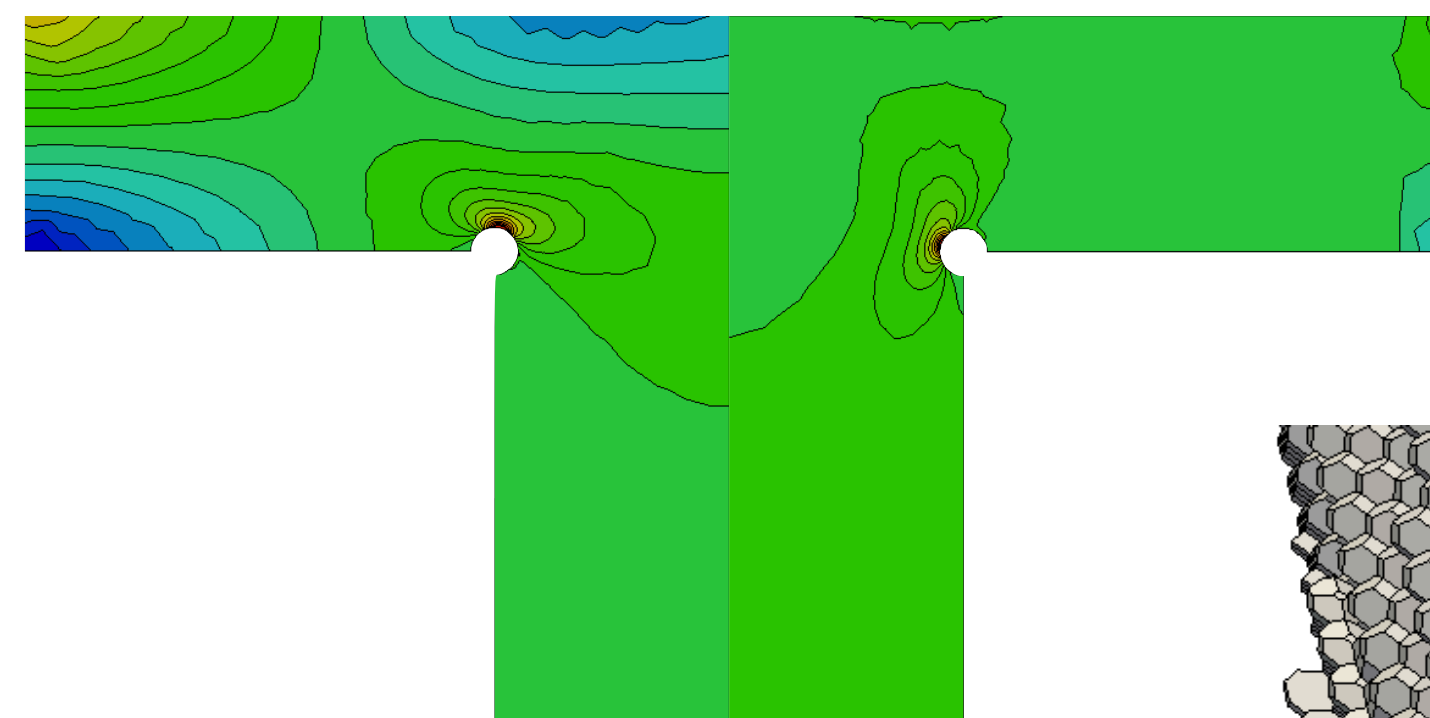
(a) Demirdžić et al. [3] Benchmark



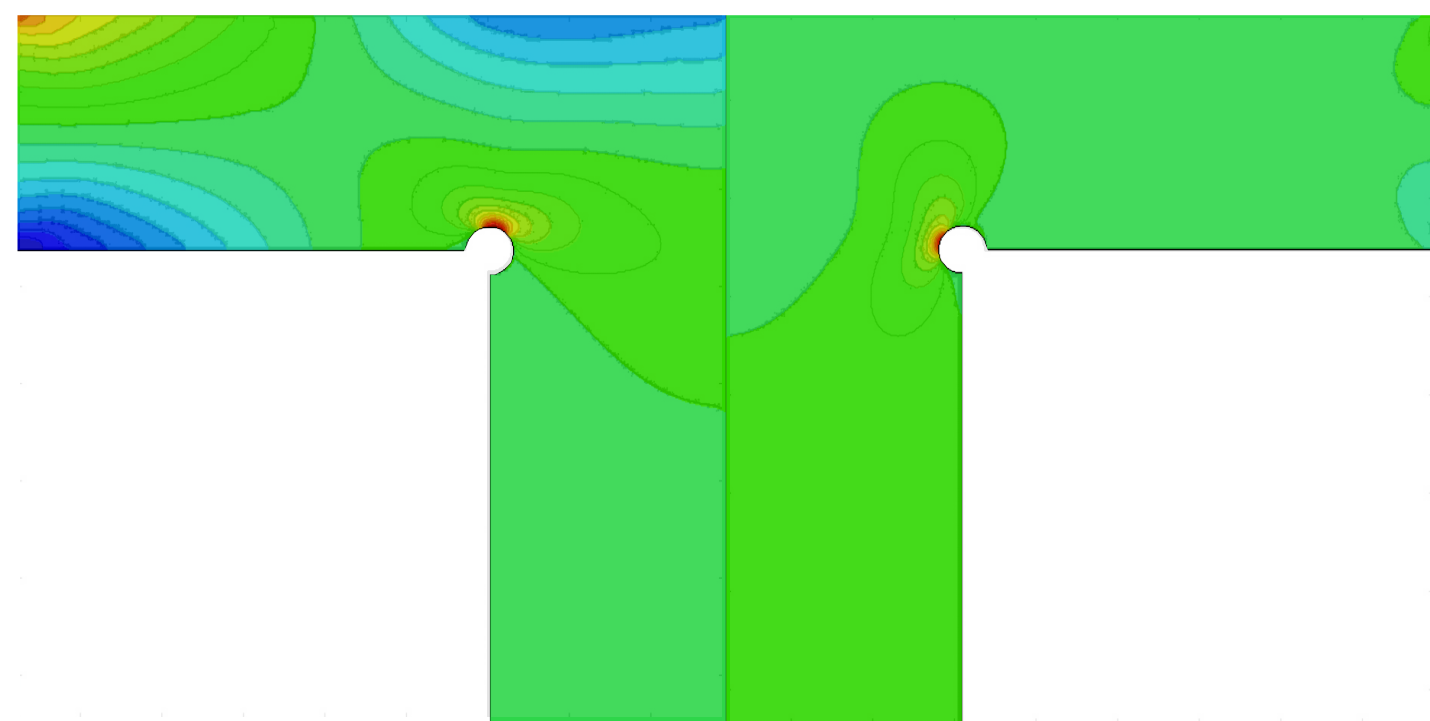
(b) Hexahedral Mesh



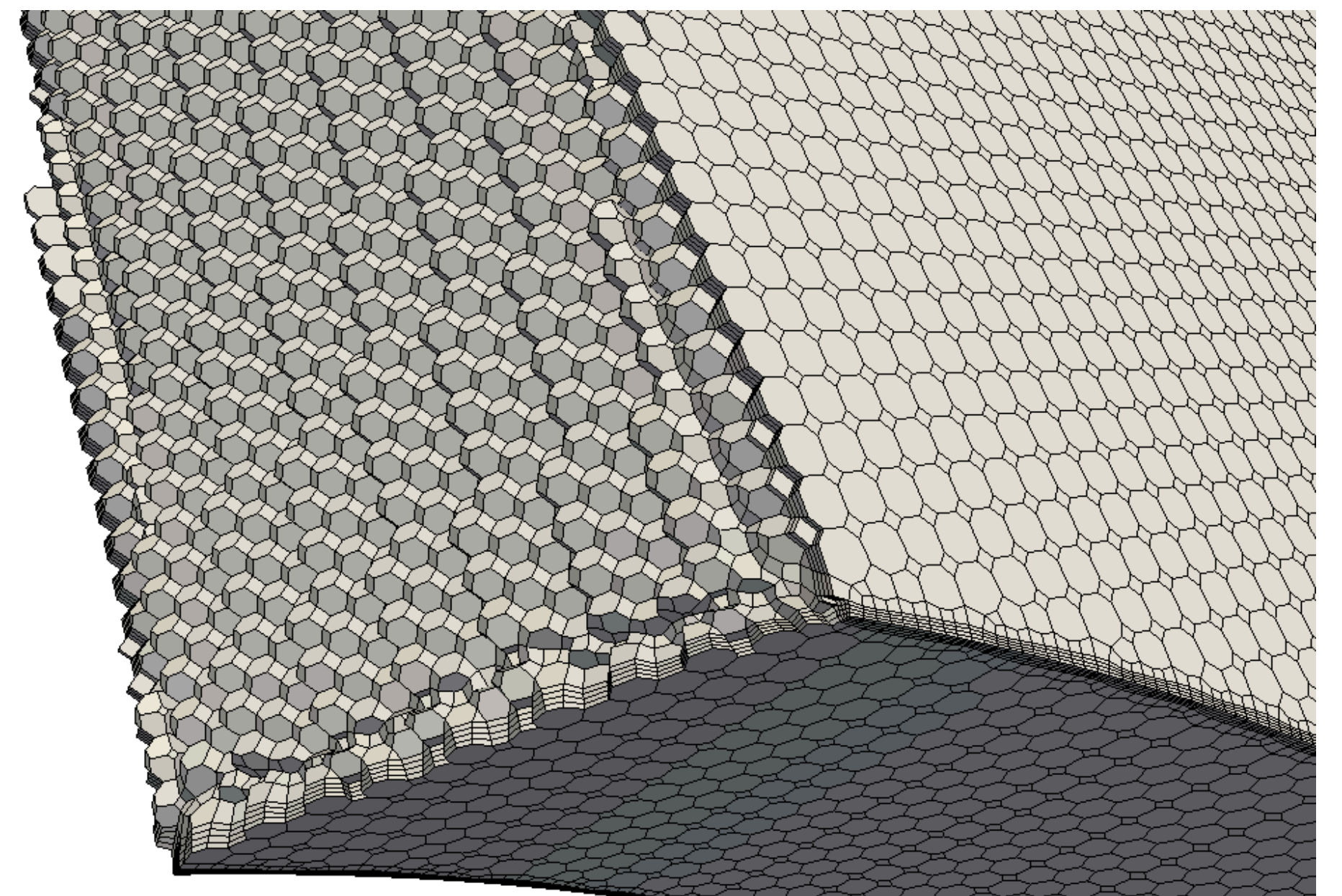
(c) Polyhedral Mesh



(d) Tetrahedral Mesh



(e) Abaqus Hexahedral Mesh



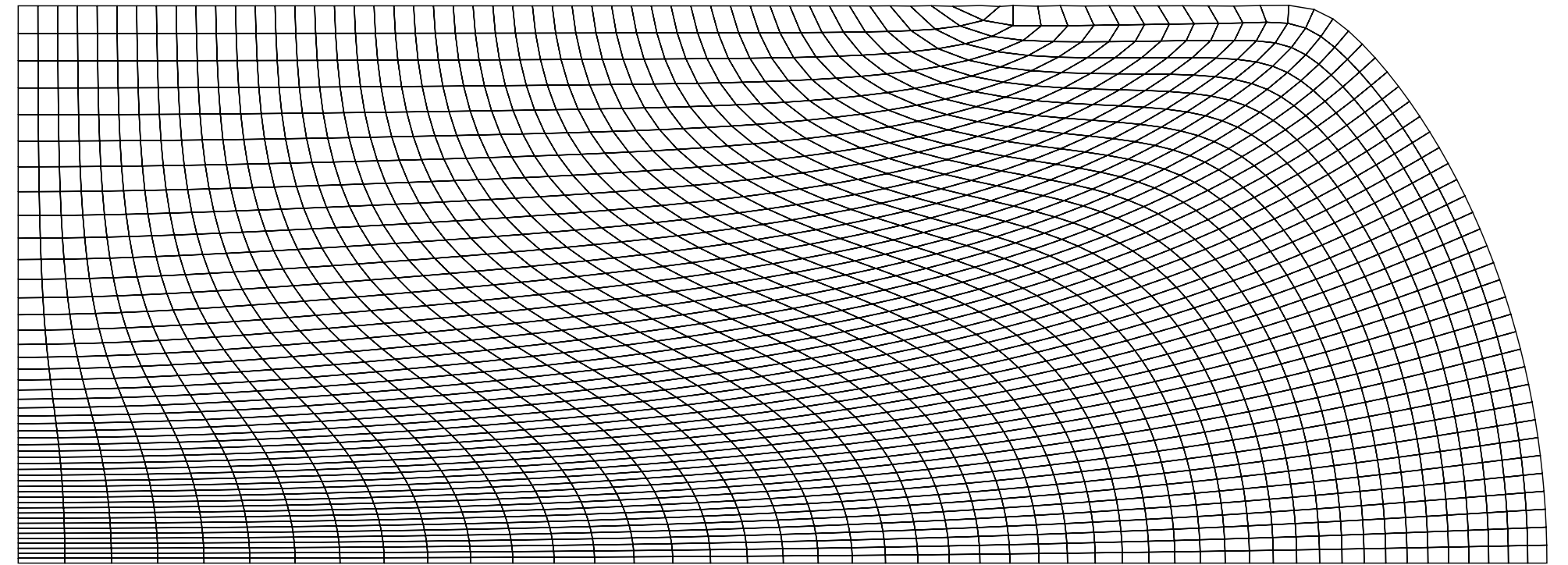
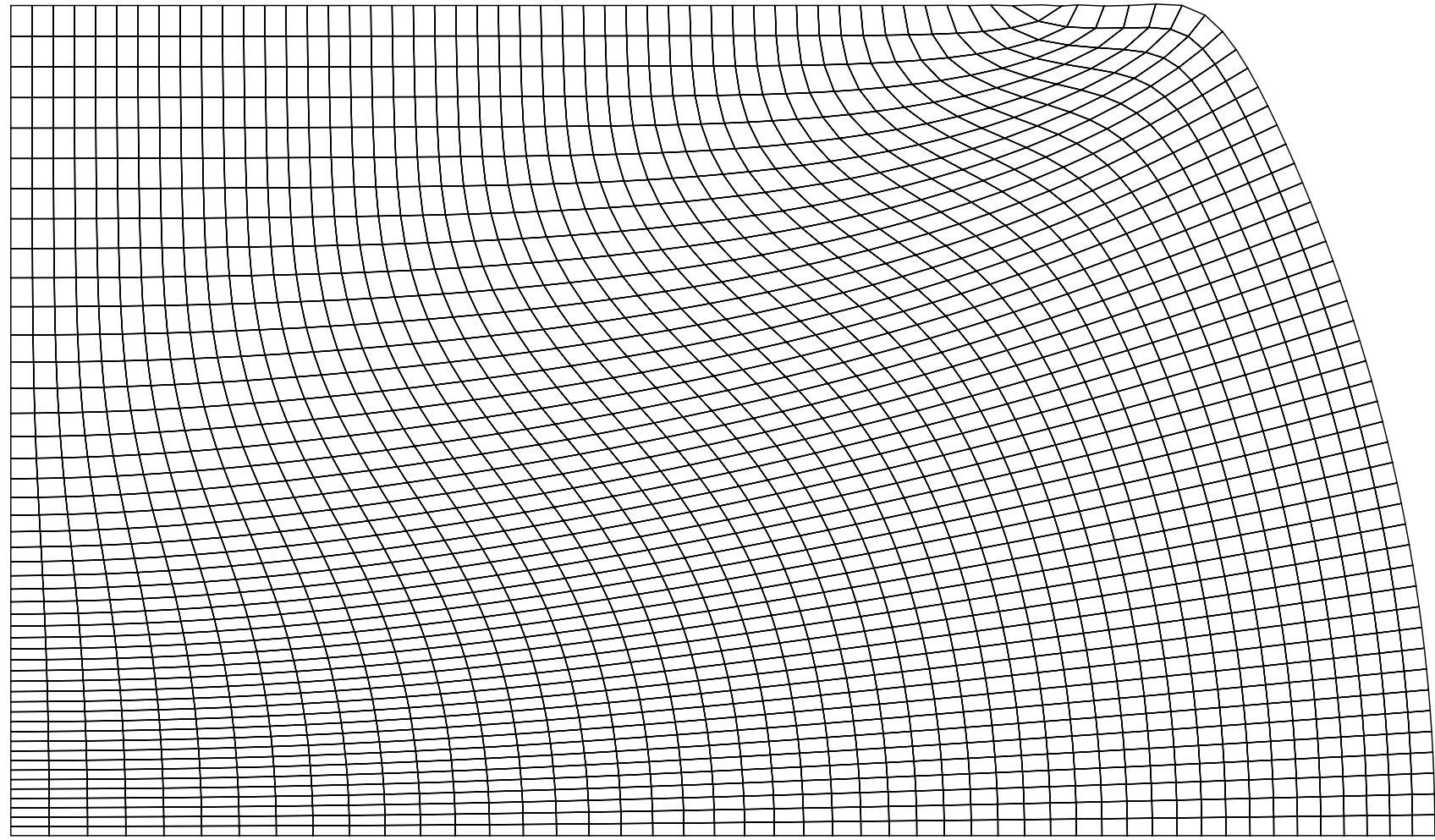
# Narrow T-section Component

Mesh	Coupled		Segregated <sup>1</sup>		Segregated <sup>2</sup>		Abaqus	
	Time	Memory	Time	Memory	Time	Memory	Time	Memory
624	0.2	15	0.7	8	1.5	10	2	98
4 992	2	58	6	24	10	30	3	102
39 936	29	340	98	88	416	140	33	1 500
319 488	421	2 400	2 220	560	2 647	900	1 236	13 000

Table III. Narrow T-section component: wall-clock time (in s) and maximum memory usage (in MB)

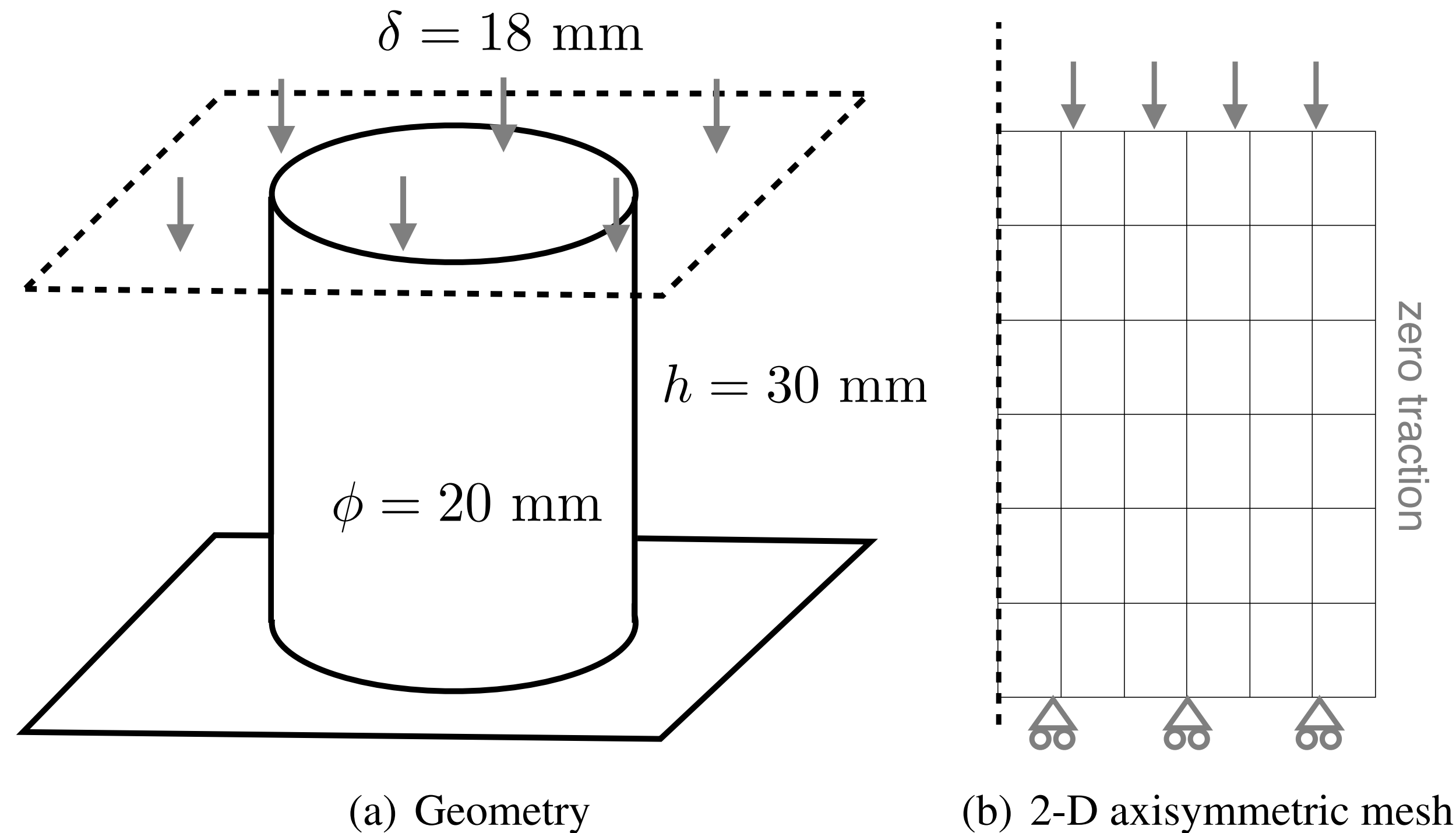
*See article almost in press \*fingers crossed\*:*

P. Cardiff, Z. Tukovic, H. Jasak, and A. Ivankovic, “**A block-coupled finite volume methodology for linear elasticity and unstructured meshes,**” *Computers and structures*, 2016, *revision under review*.



# Nonlinear Benchmark Cases

# Upsetting a billet



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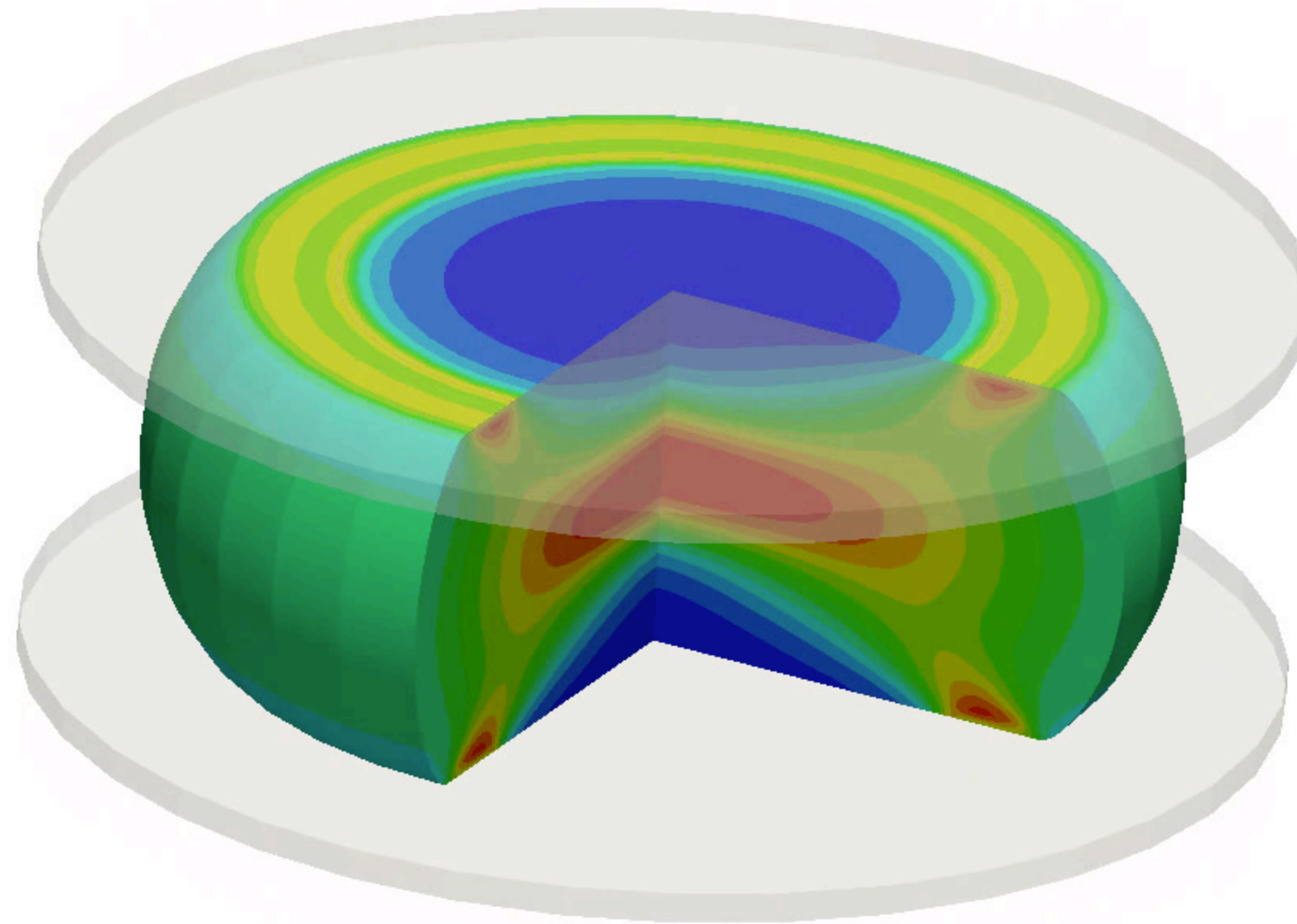
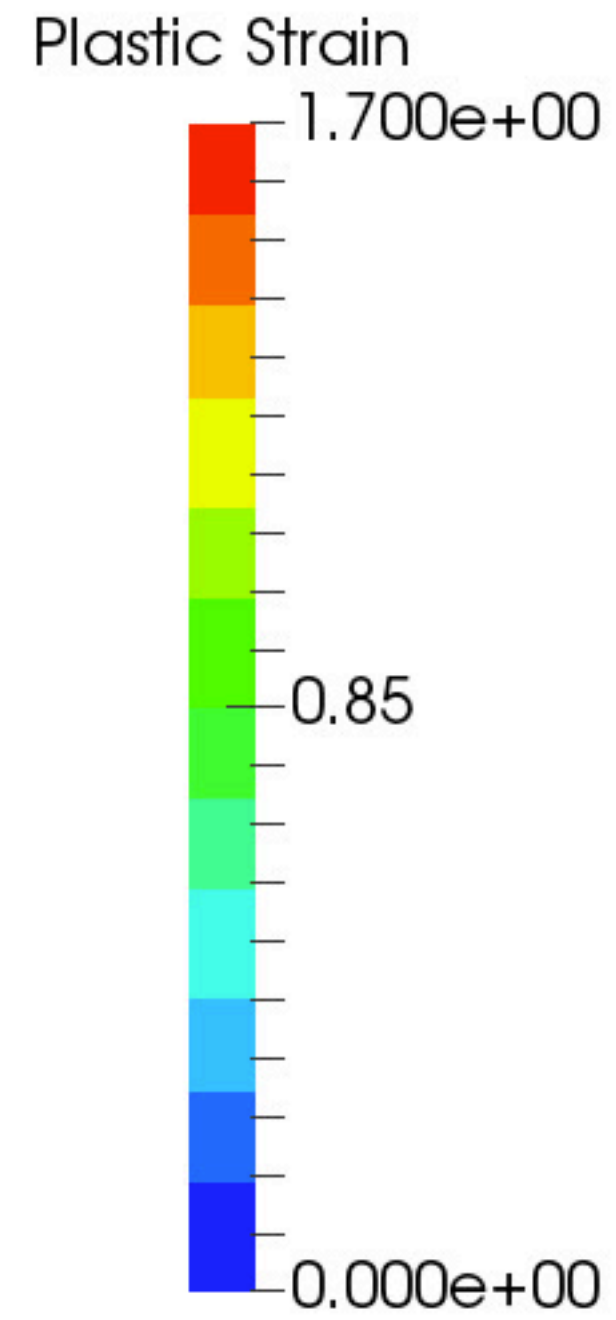
Young's modulus	$E$	200 GPa
Poisson's ratio	$\nu$	0.3
Initial yield stress	$\sigma_Y$	700 MPa
Hardening parameter	$\kappa$	300 MPa

---

See *other* article almost in press \*fingers crossed\*:

P. Cardiff, Z. Tukovic, P. De Jaeger, M. Clancy and A. Ivankovic, “**A Lagrangian cell-centred finite volume method for metal forming simulation**,” International journal for numerical methods in engineering, 2016, *revision under review*.

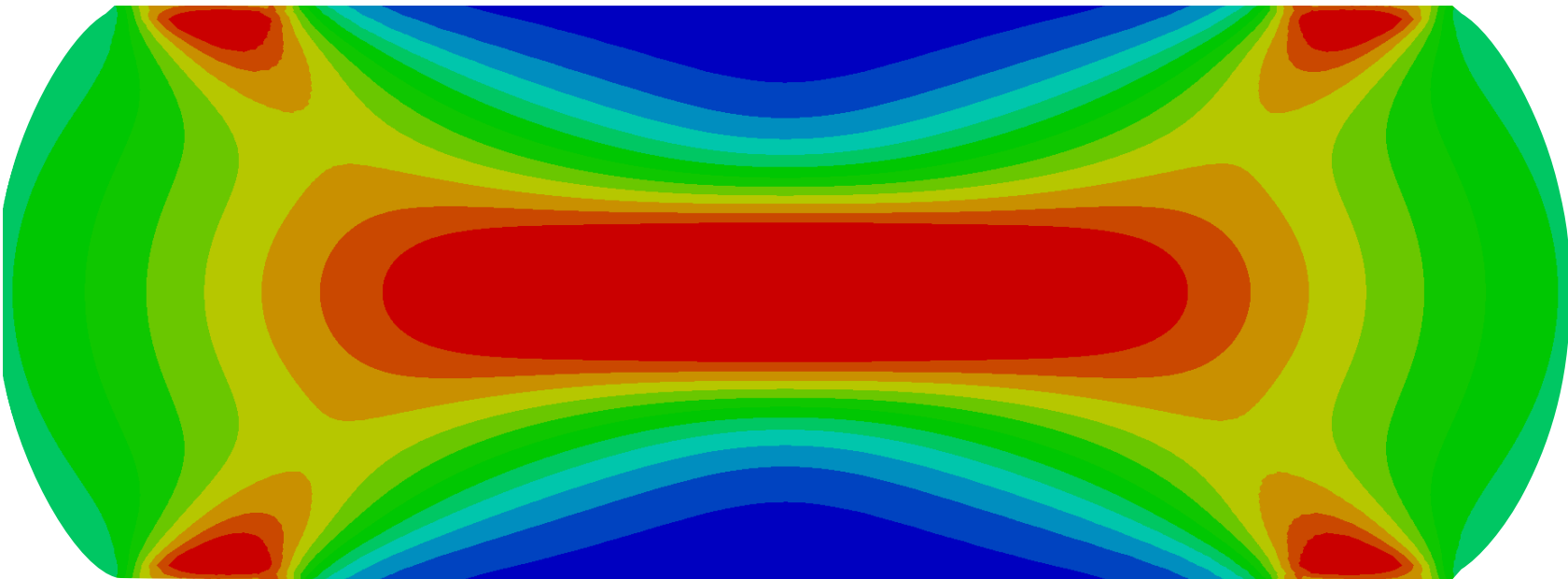
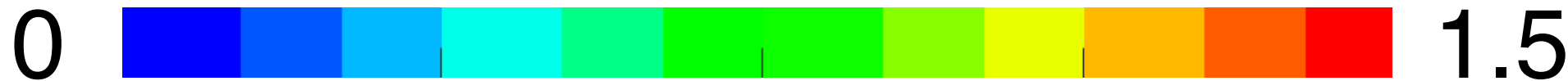
# Upsetting a billet



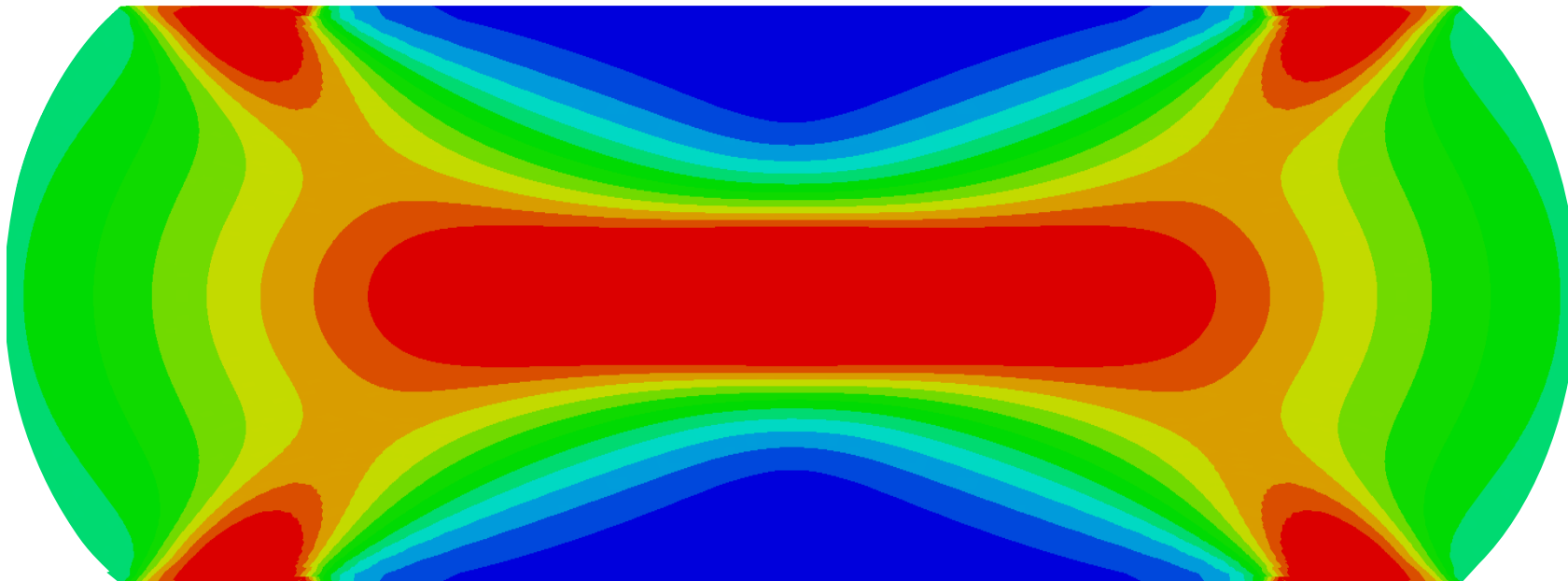
Precent upset: 60.0

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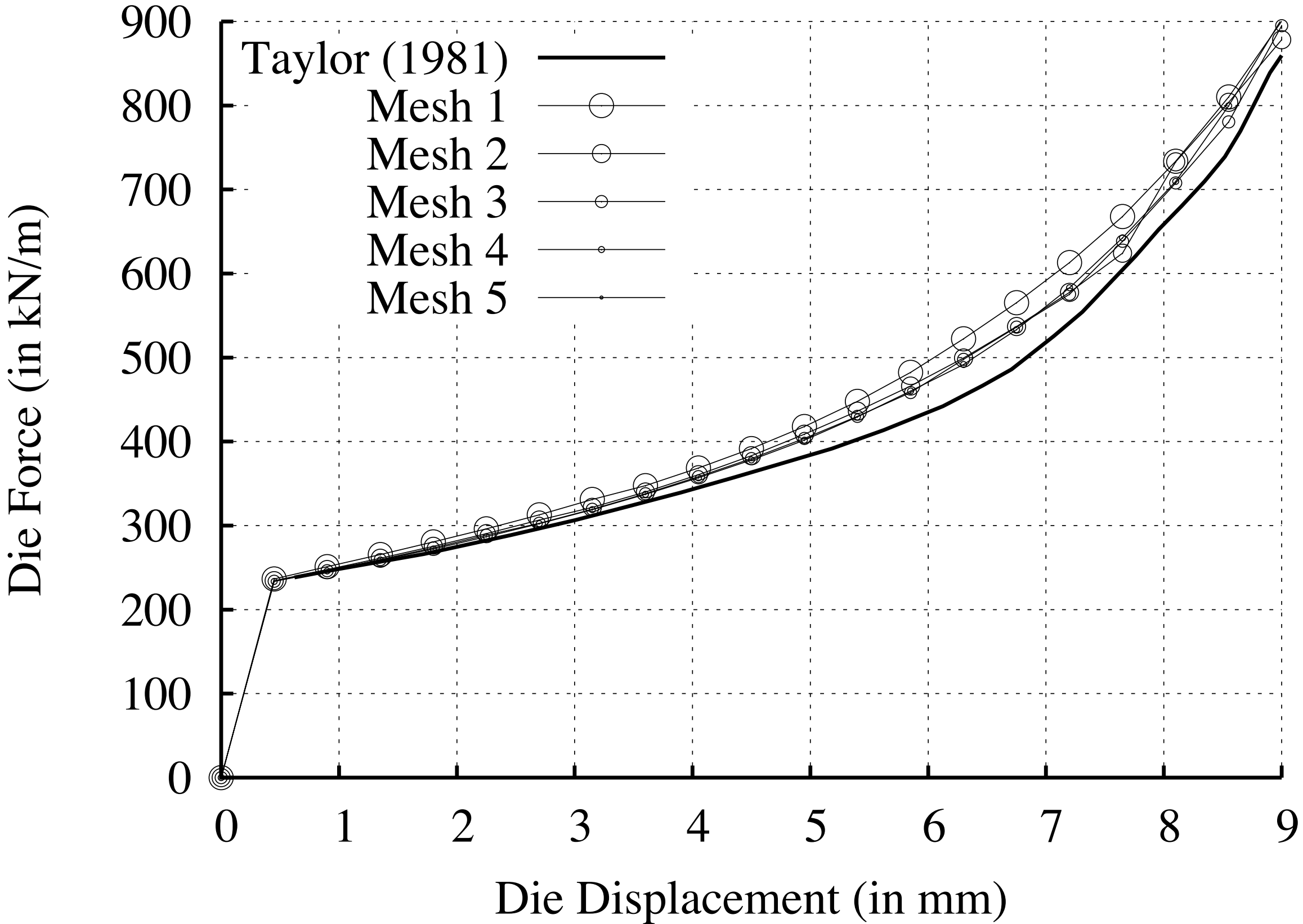
# Upsetting a billet



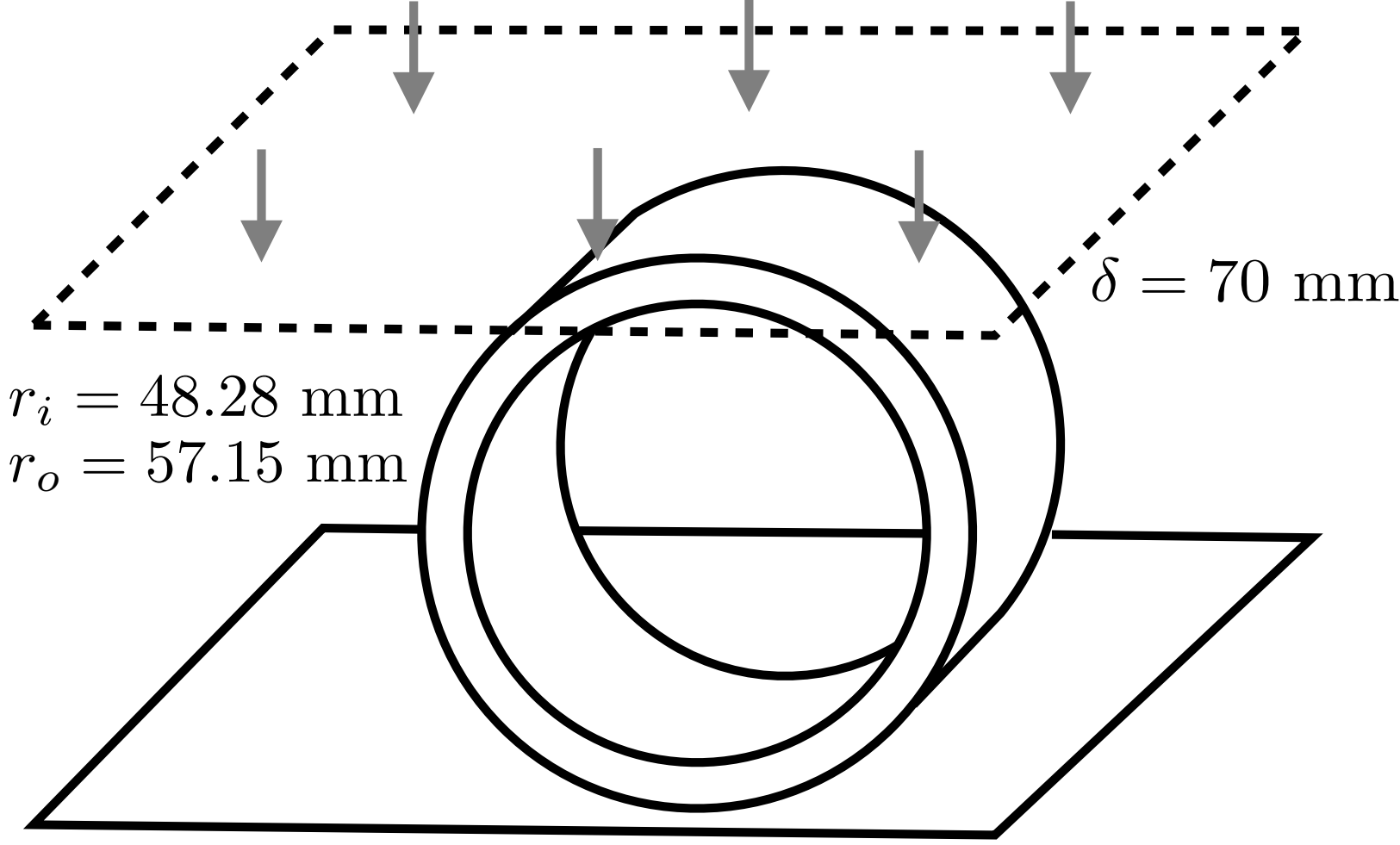
OpenFOAM vertex-wise



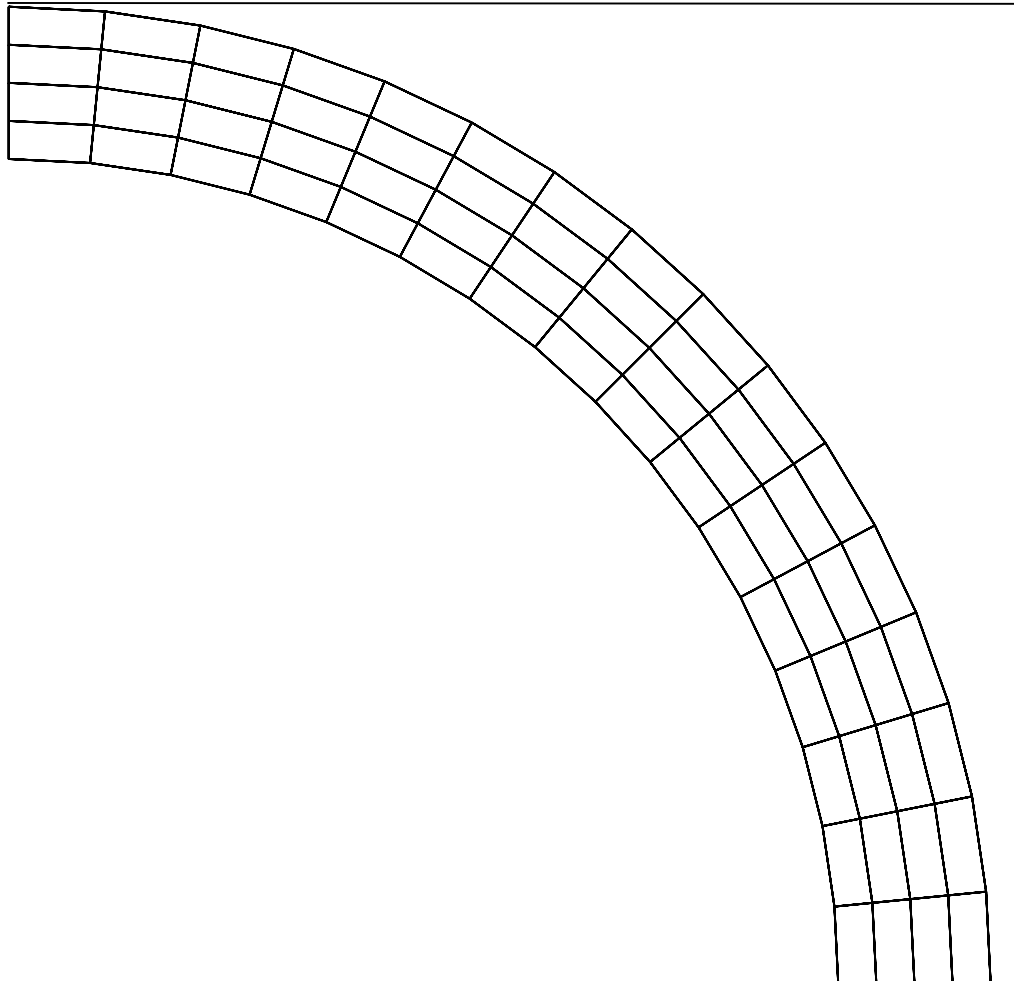
Abaqus



# Crushing a pipe



(a) Geometry



(b) 2-D Plane Strain Mesh

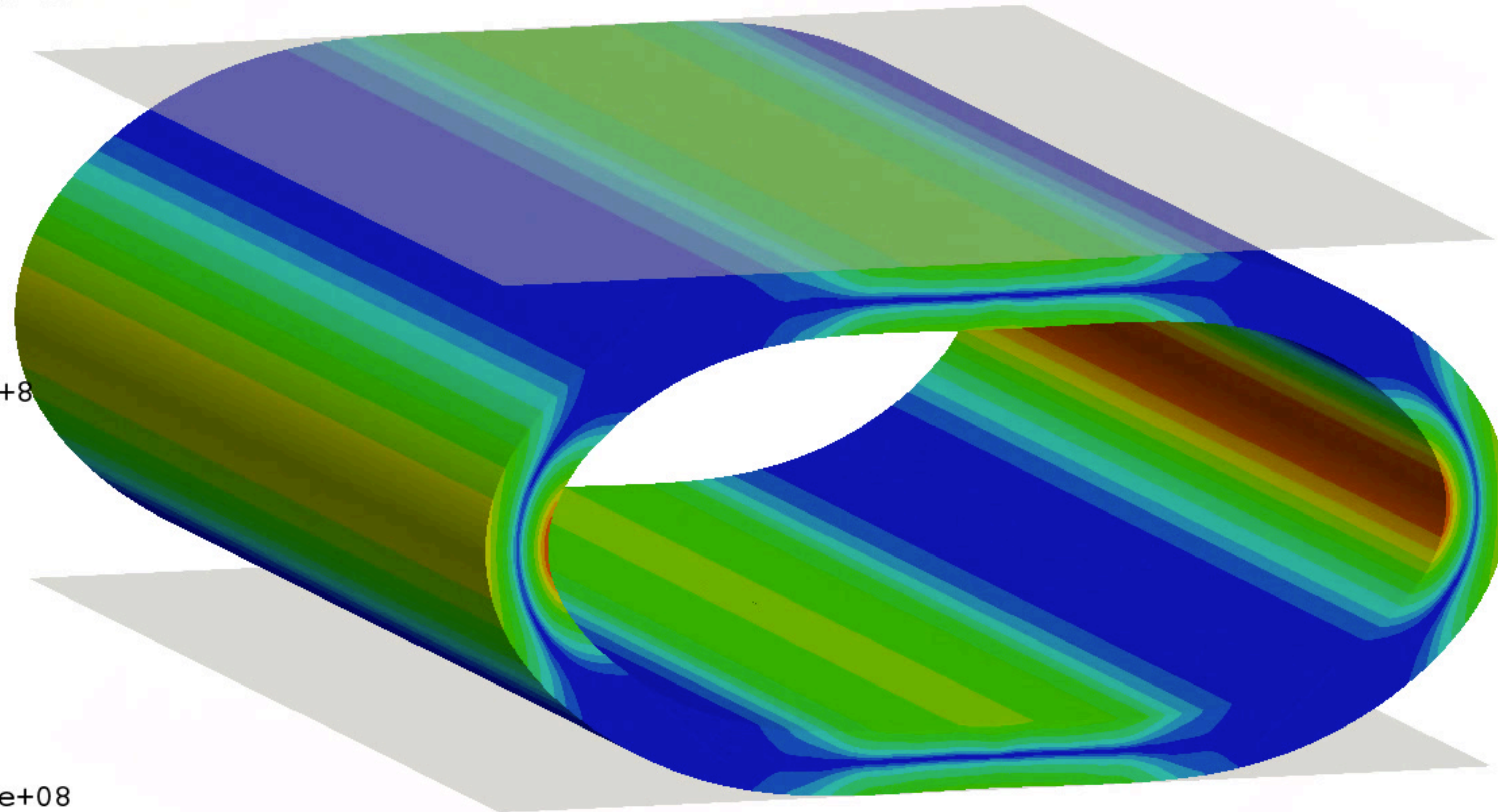
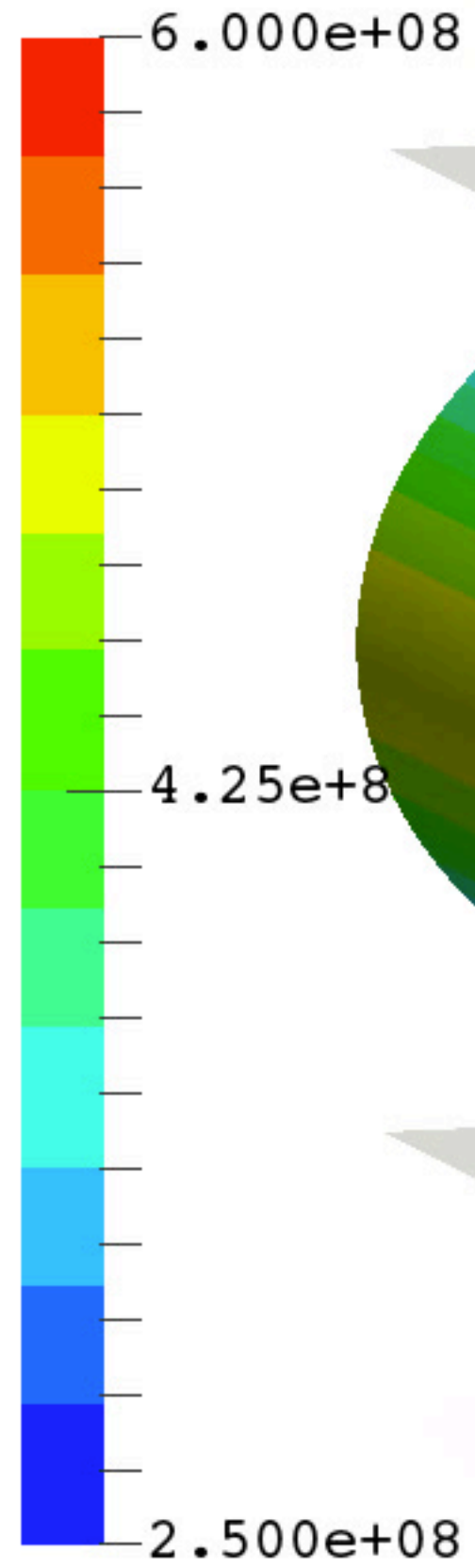
---

Young's modulus	$E$	186 GPa
Poisson's ratio	$\nu$	0.3
Initial yield stress	$\sigma_Y$	241.32 MPa
Hardening	Plastic strain	Yield stress (in MPa)
	0	241.32
	0.0035	275.79
	0.0083	301.65
	0.0133	318.88
	0.0182	344.74
	0.0281	361.98
	0.0380	379.21
	1	1 482.37

---

# Crushing of a pipe

Strength (in Pa)

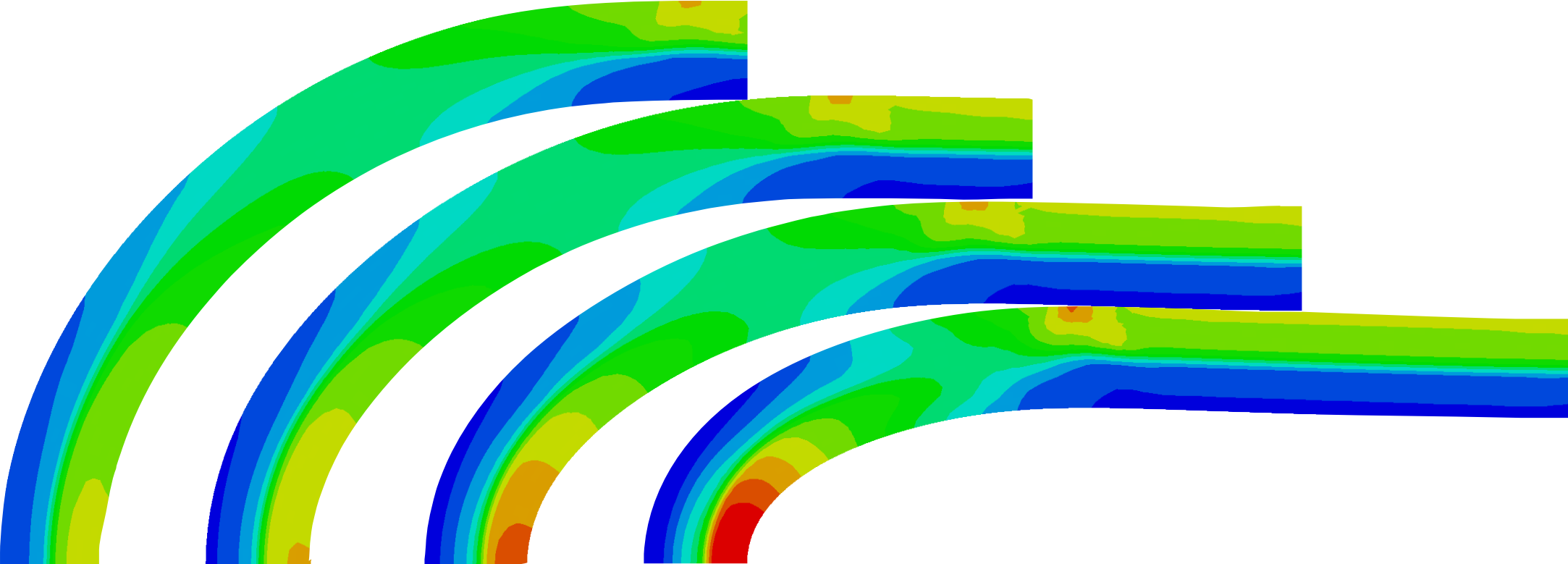


Time: 0.61 s

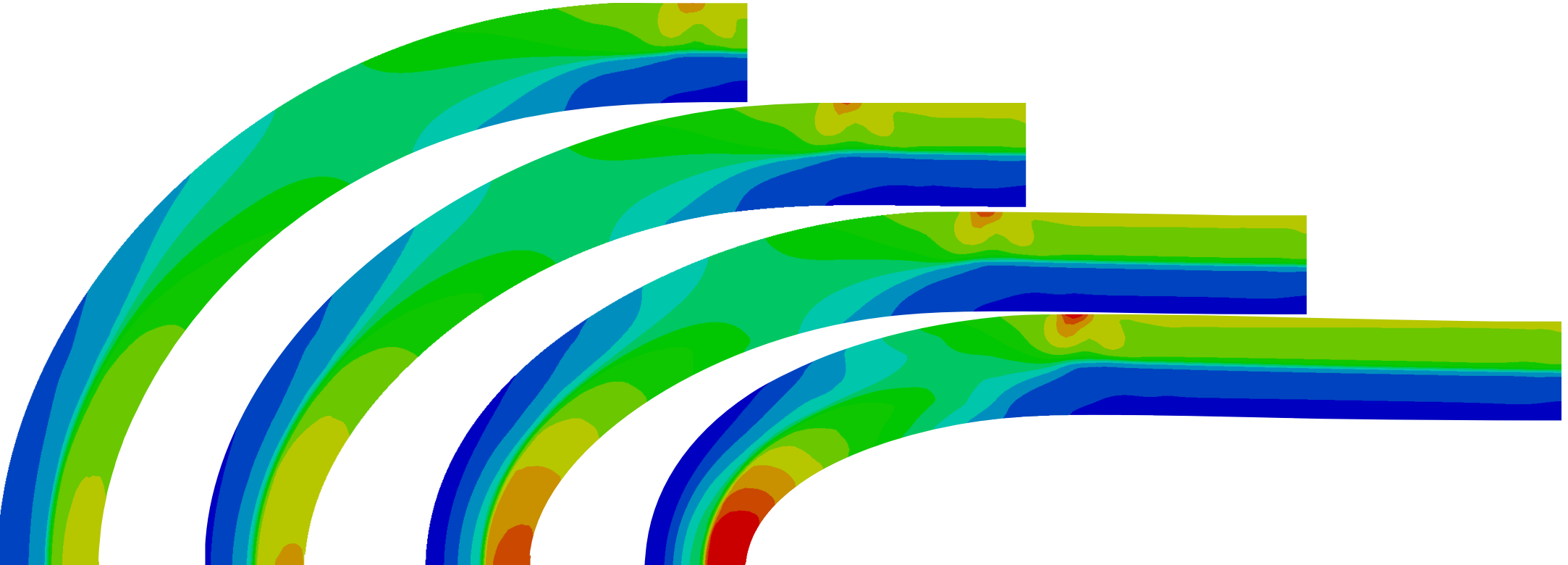
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# Crushing a pipe

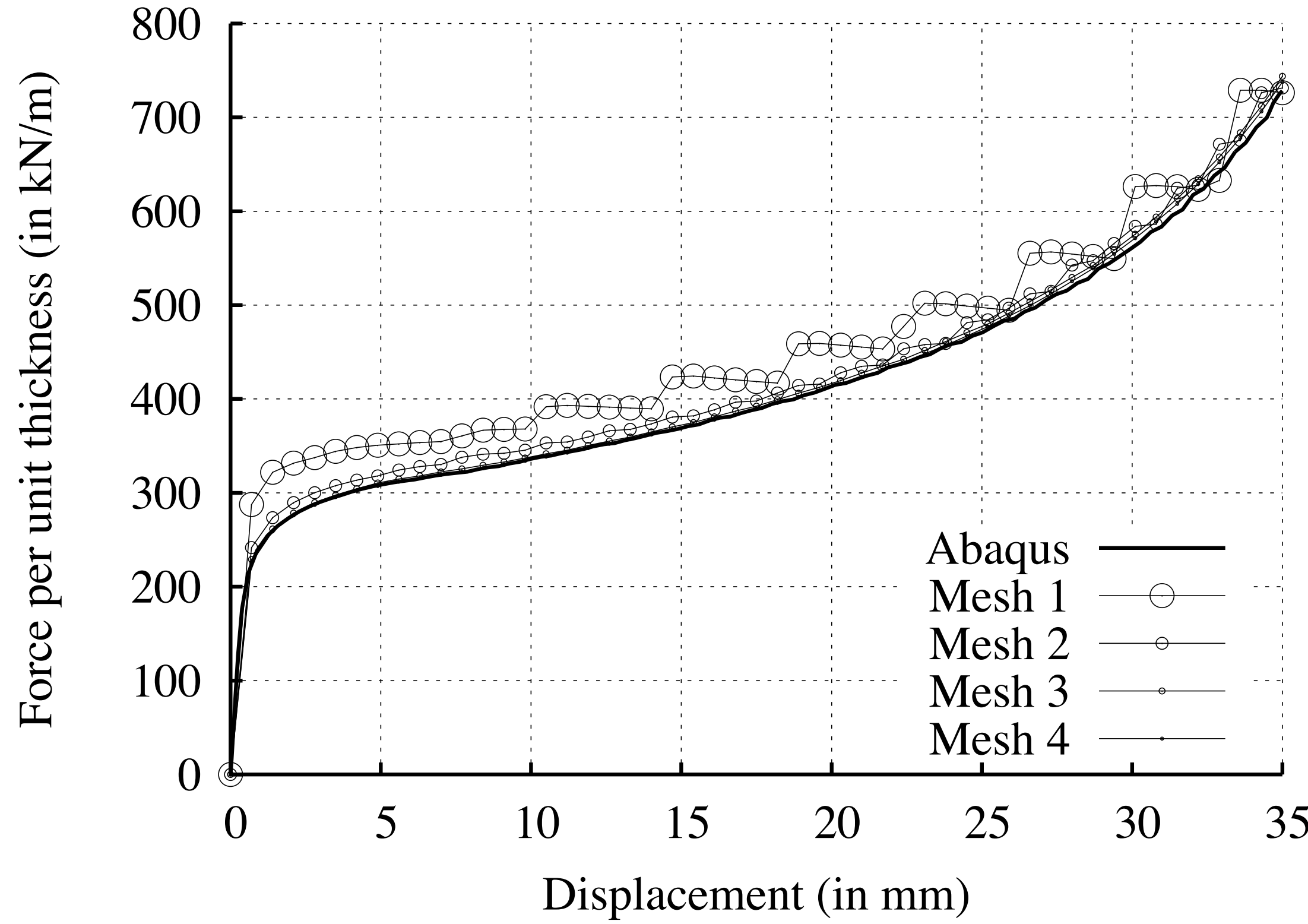
-0.3 GPa  0.5 GPa



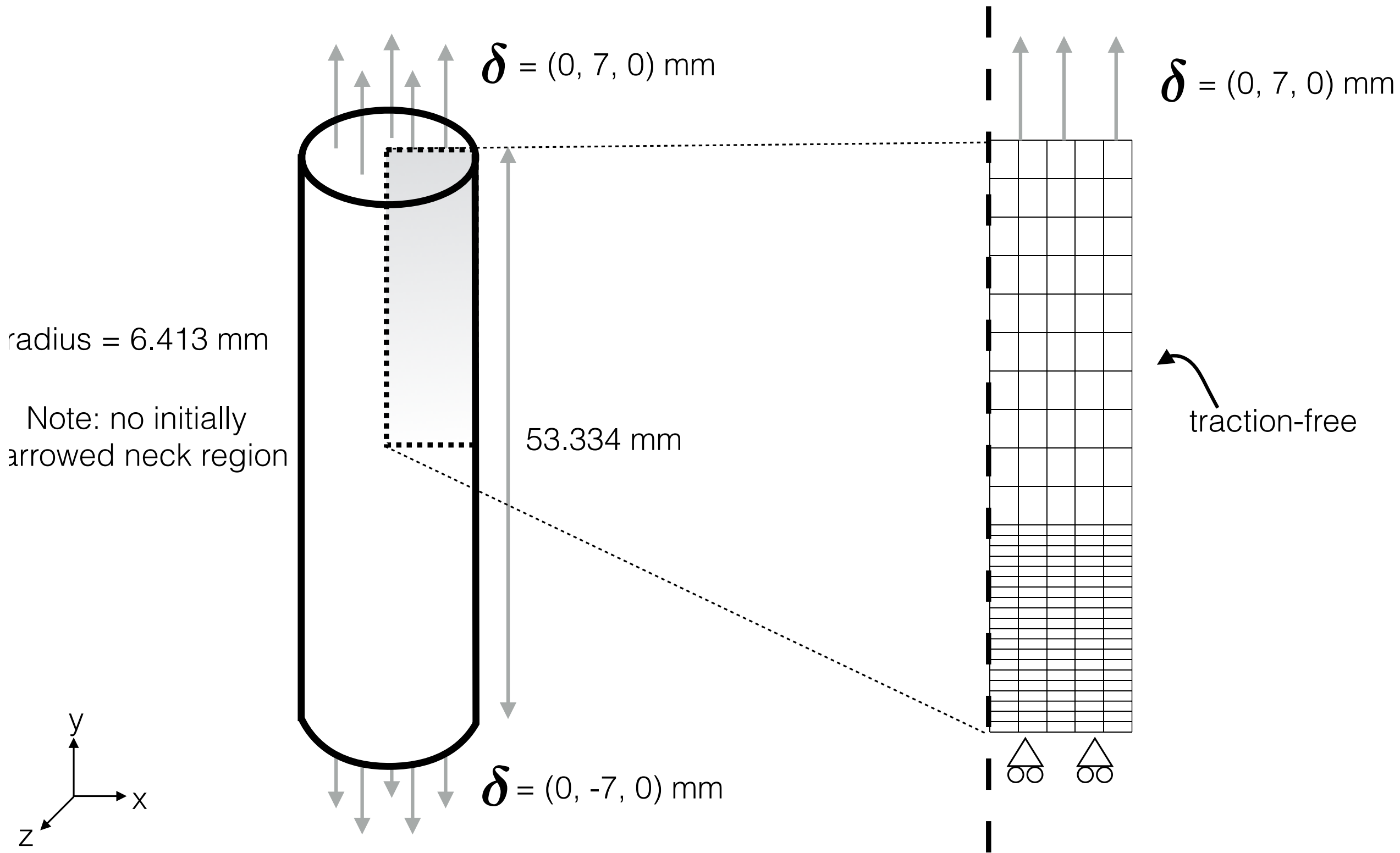
Abaqus



OpenFOAM vertex-wise (modified approach)

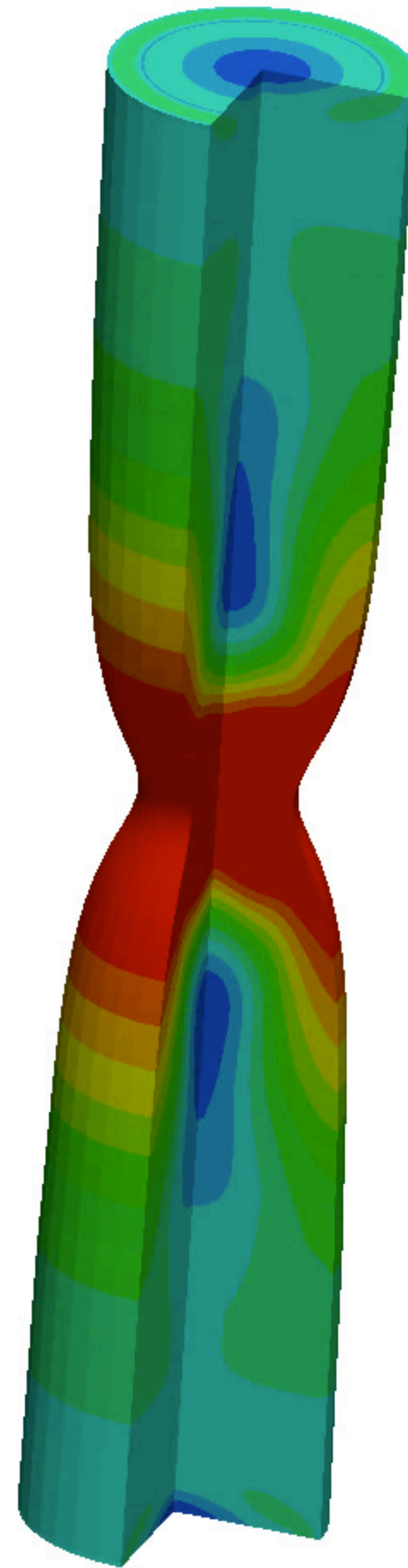
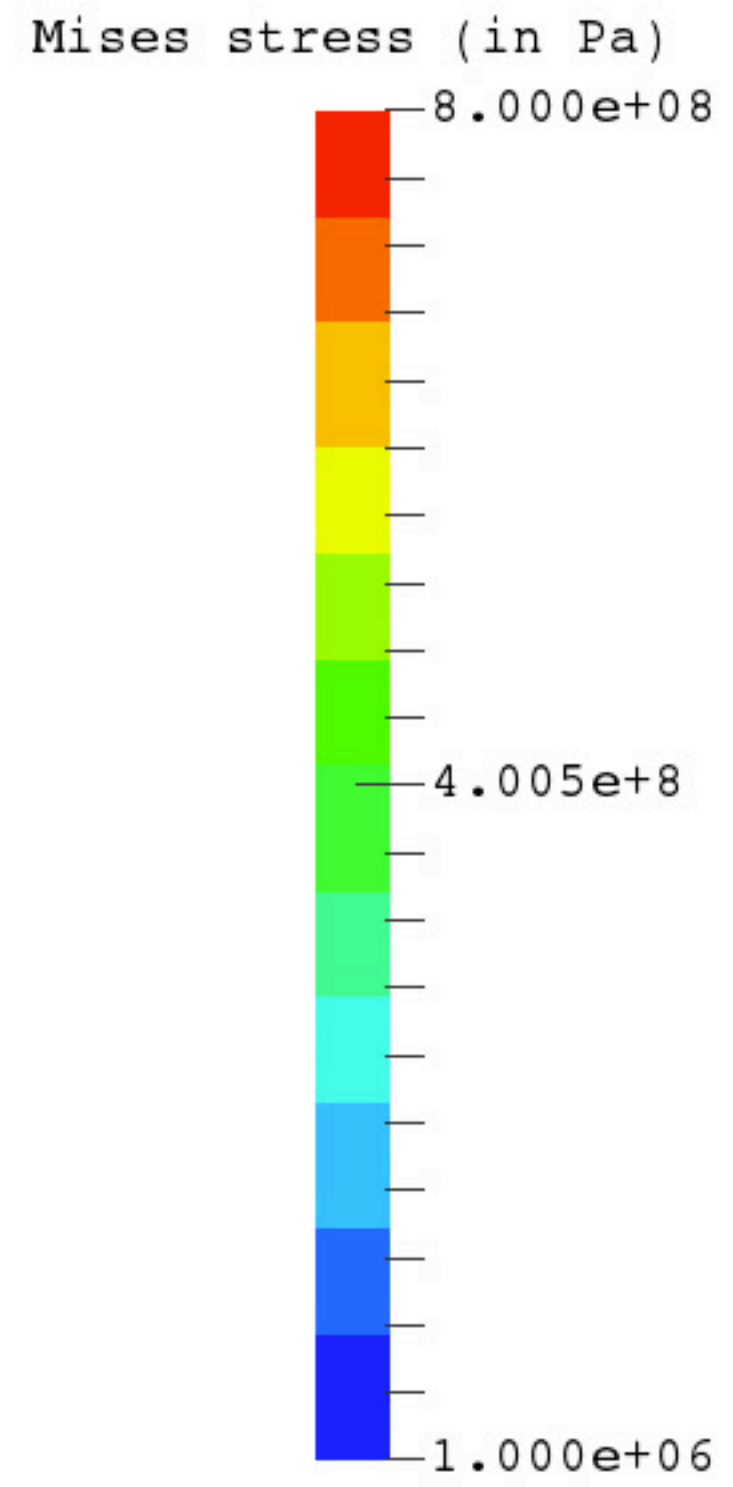


# Necking of a bar



Young's modulus	$E$	200 GPa
Poisson's ratio	$\nu$	0.3
Initial yield stress	$\sigma_Y$	451 MPa
Hardening	Plastic strain	Yield stress (in MPa)
	0.000	451
	0.006	476
	0.019	525
	0.038	583
	0.066	642
	0.147	710
	0.500	777
	1.000	831

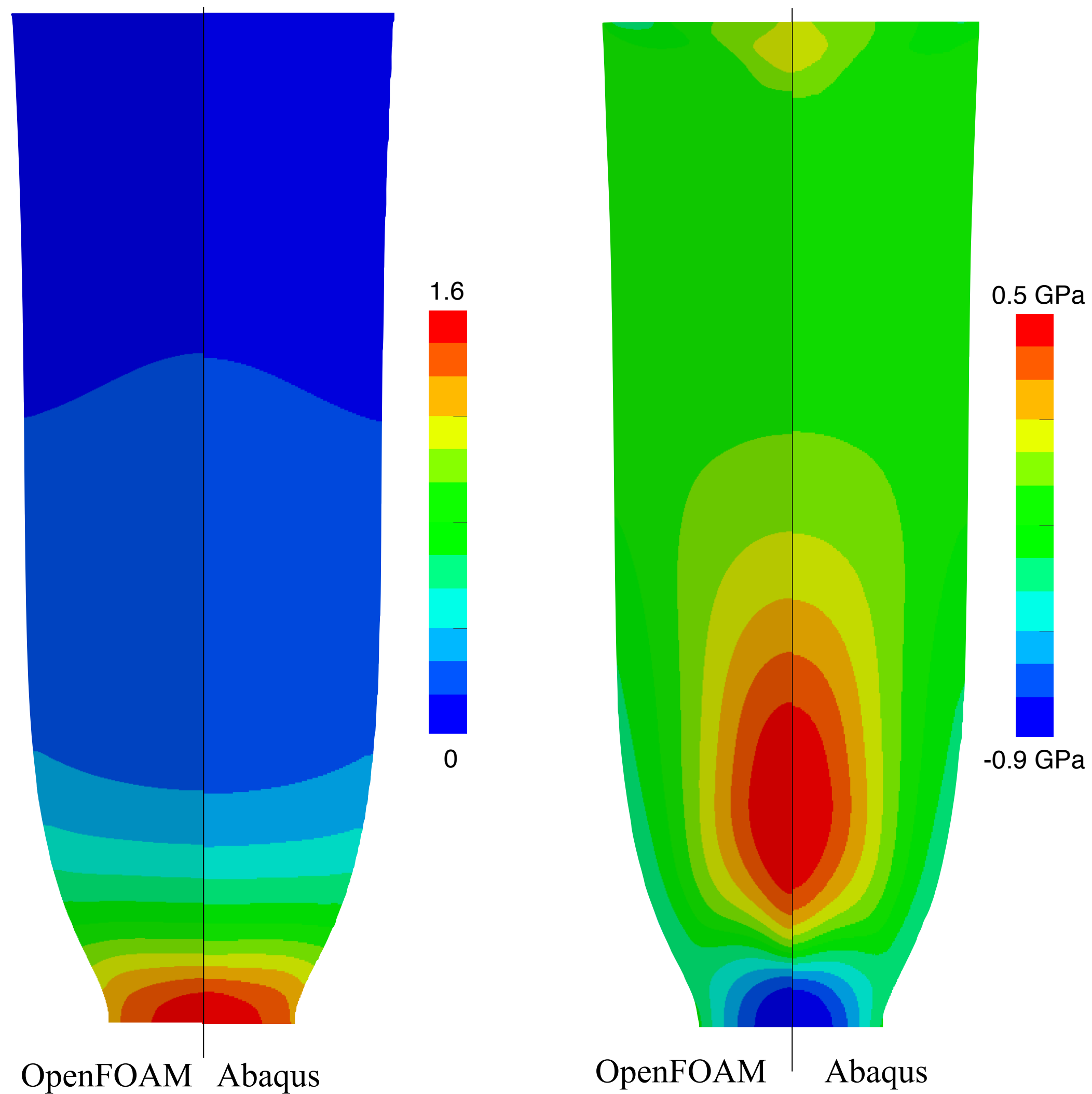
# Necking of a bar



Time: 80 s

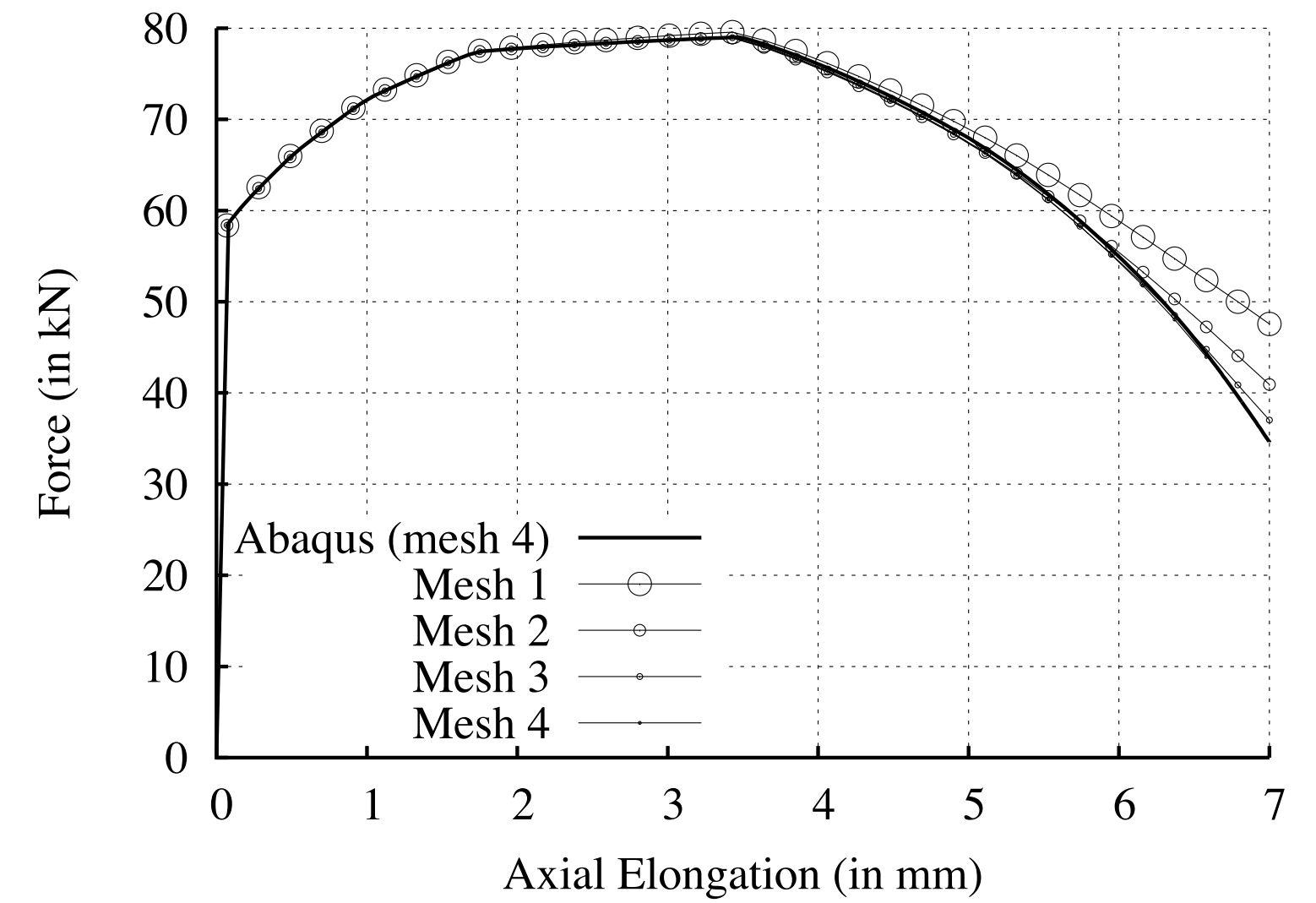
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# Necking of a bar

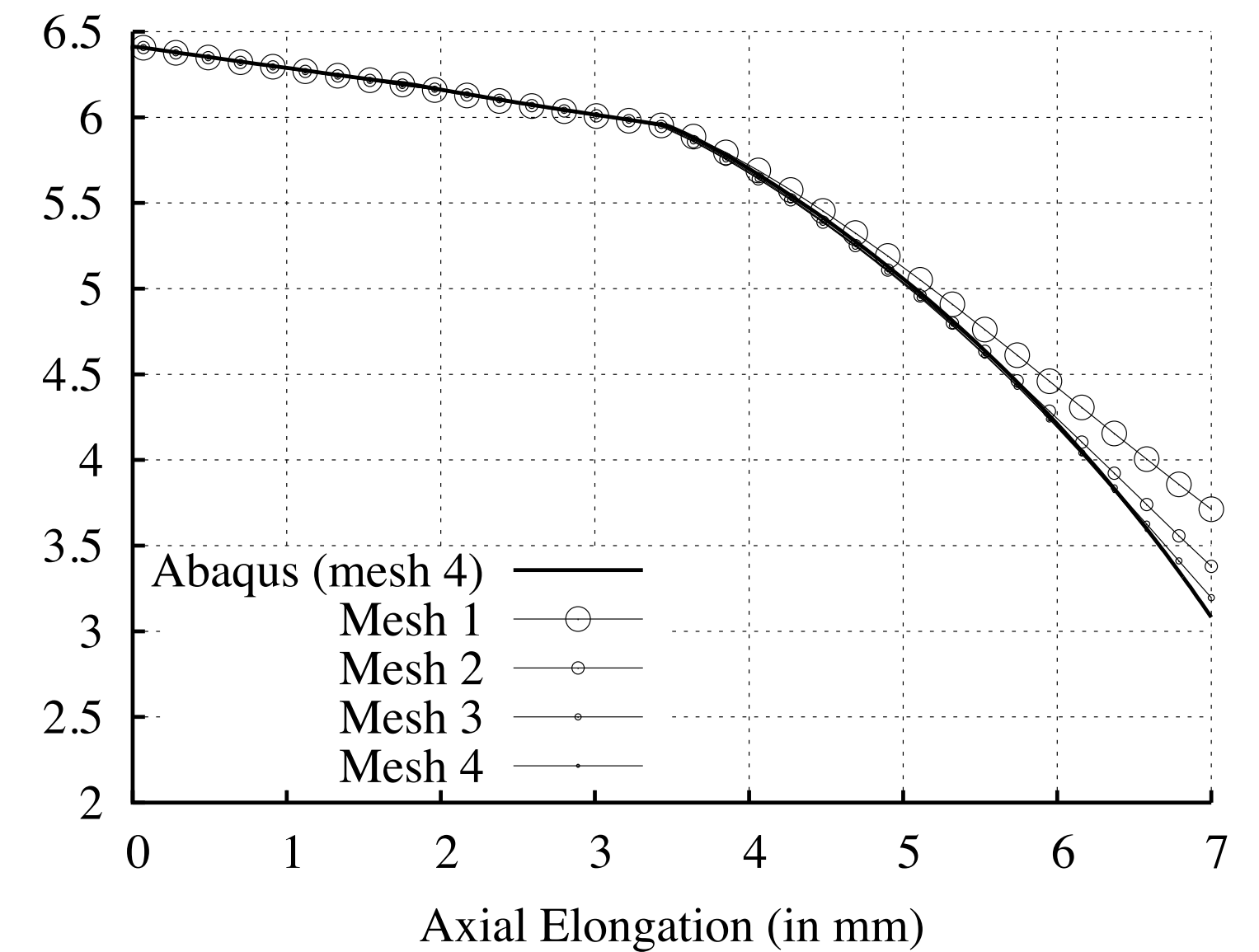


(a) Plastic strain: OpenFOAM FV (left) and Abaqus FE (right)

(b) Pressure: OpenFOAM FV (left) and Abaqus FE (right)

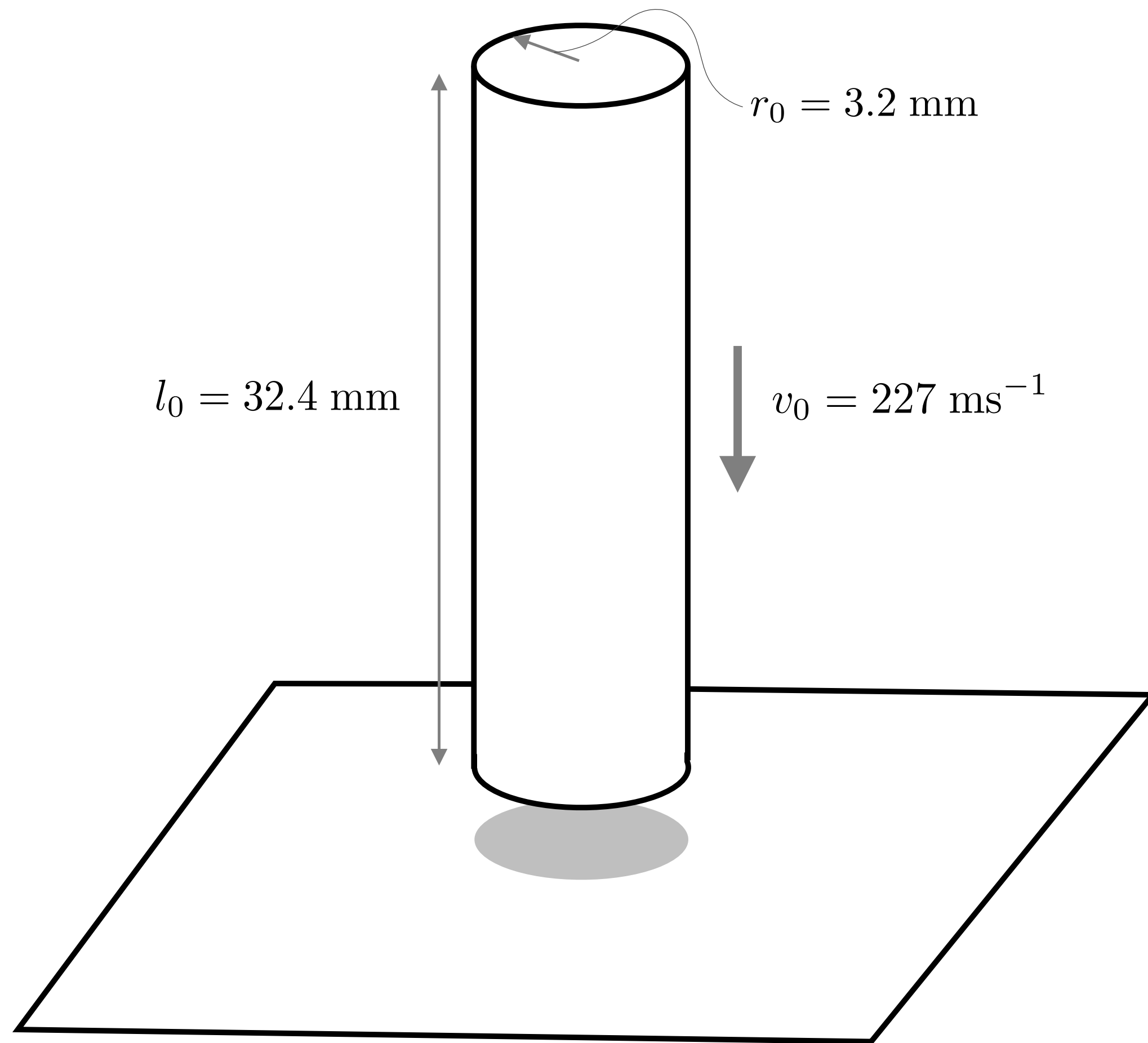


(a) Axial loading force vs. axial elongation



(b) Neck radius vs. axial elongation

# Impacting bar



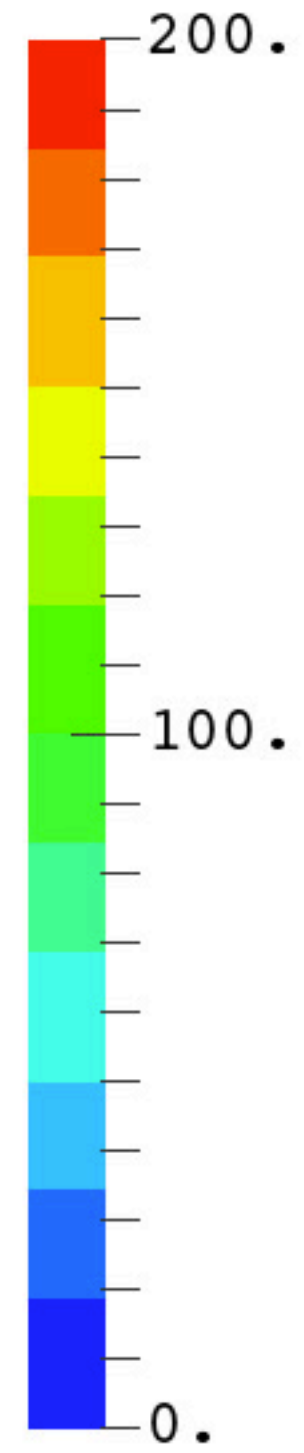
---

Initial density	$\rho$	8930 kg/m <sup>3</sup>
Young's modulus	$E$	117 GPa
Poisson's ratio	$\nu$	0.35
Initial yield stress	$\sigma_Y$	400 MPa
Hardening modulus	$\kappa$	100 MPa

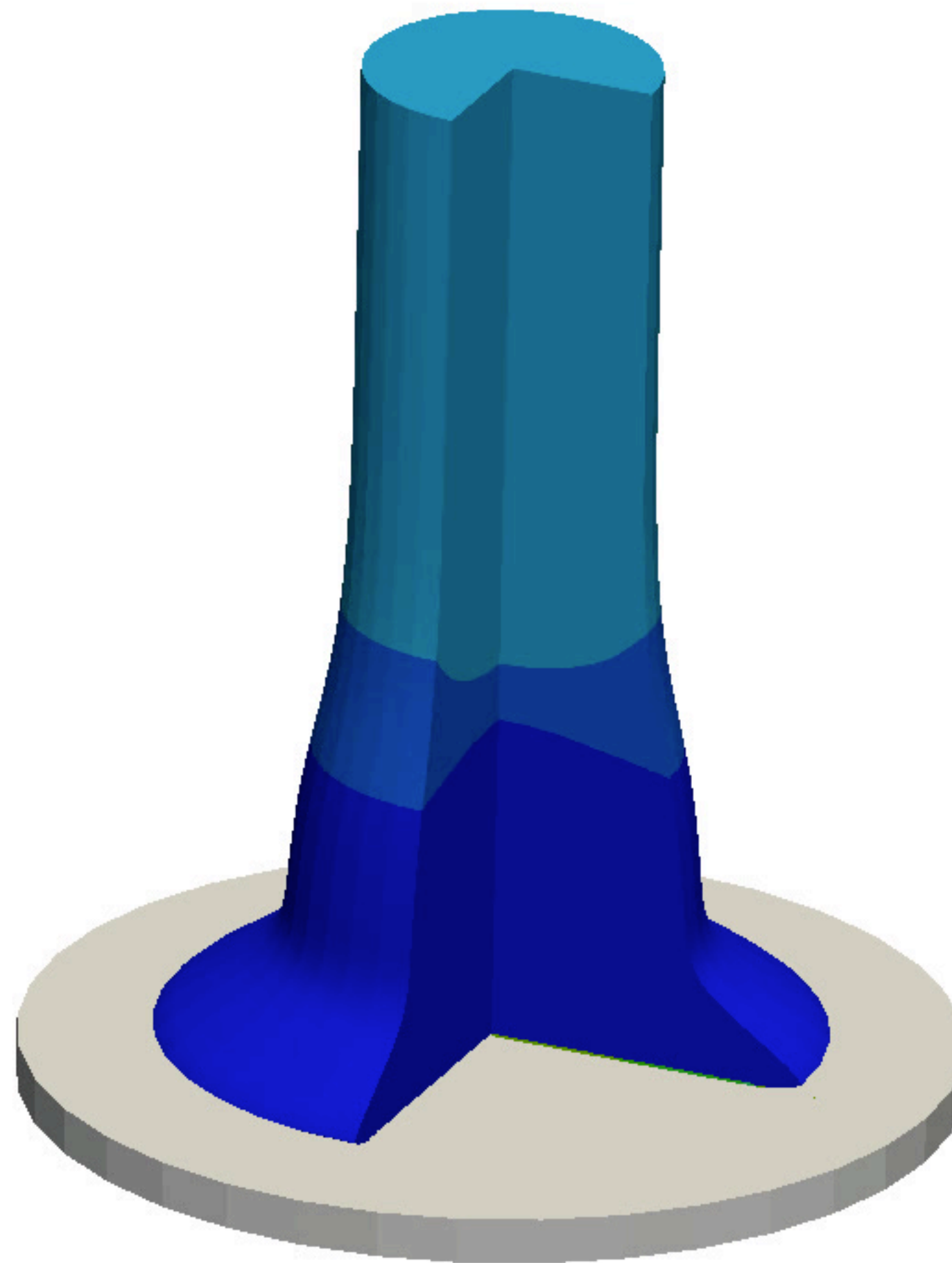
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# High speed billet impact

Velocity (in m/s)

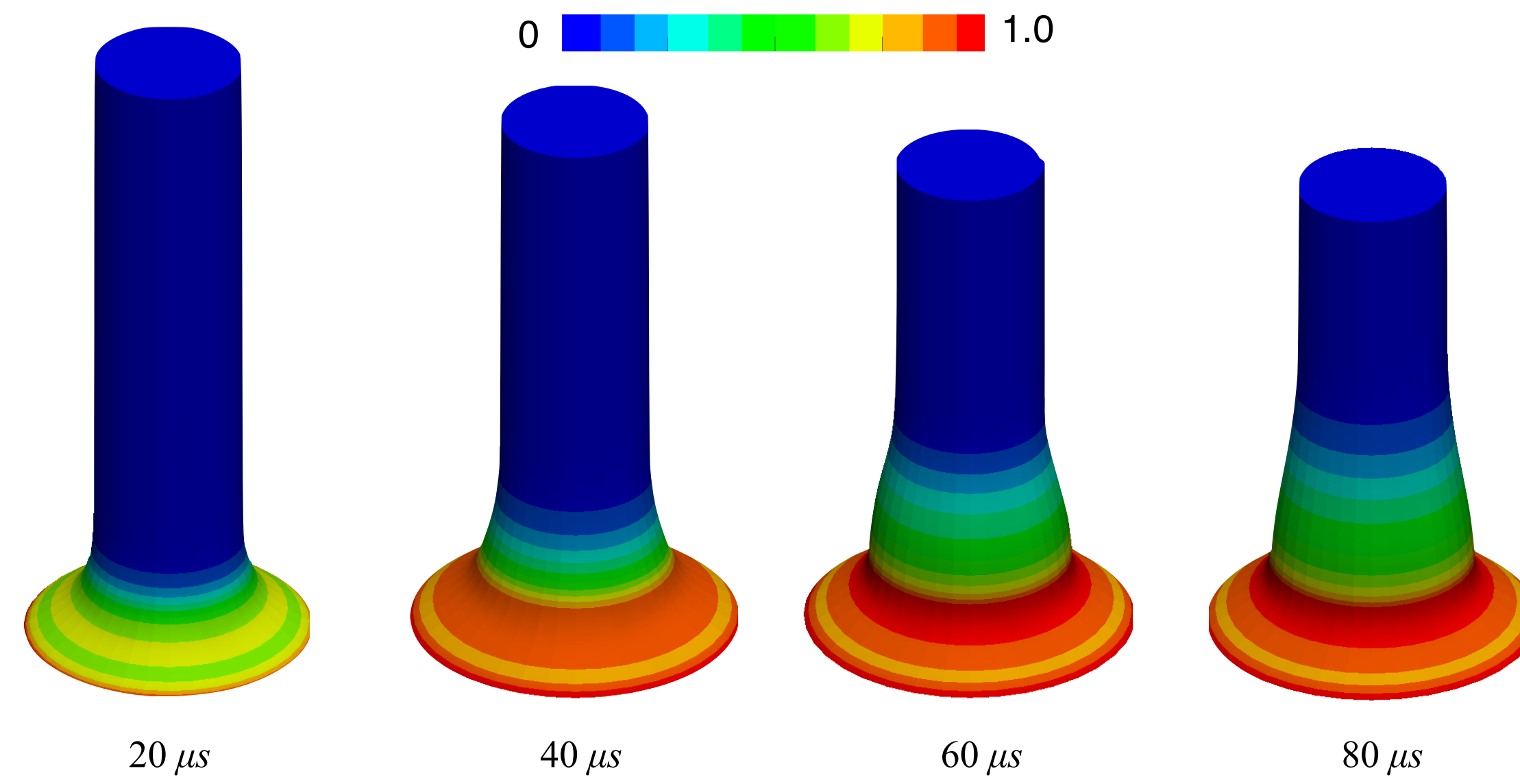


Time: 69  $\mu$ s

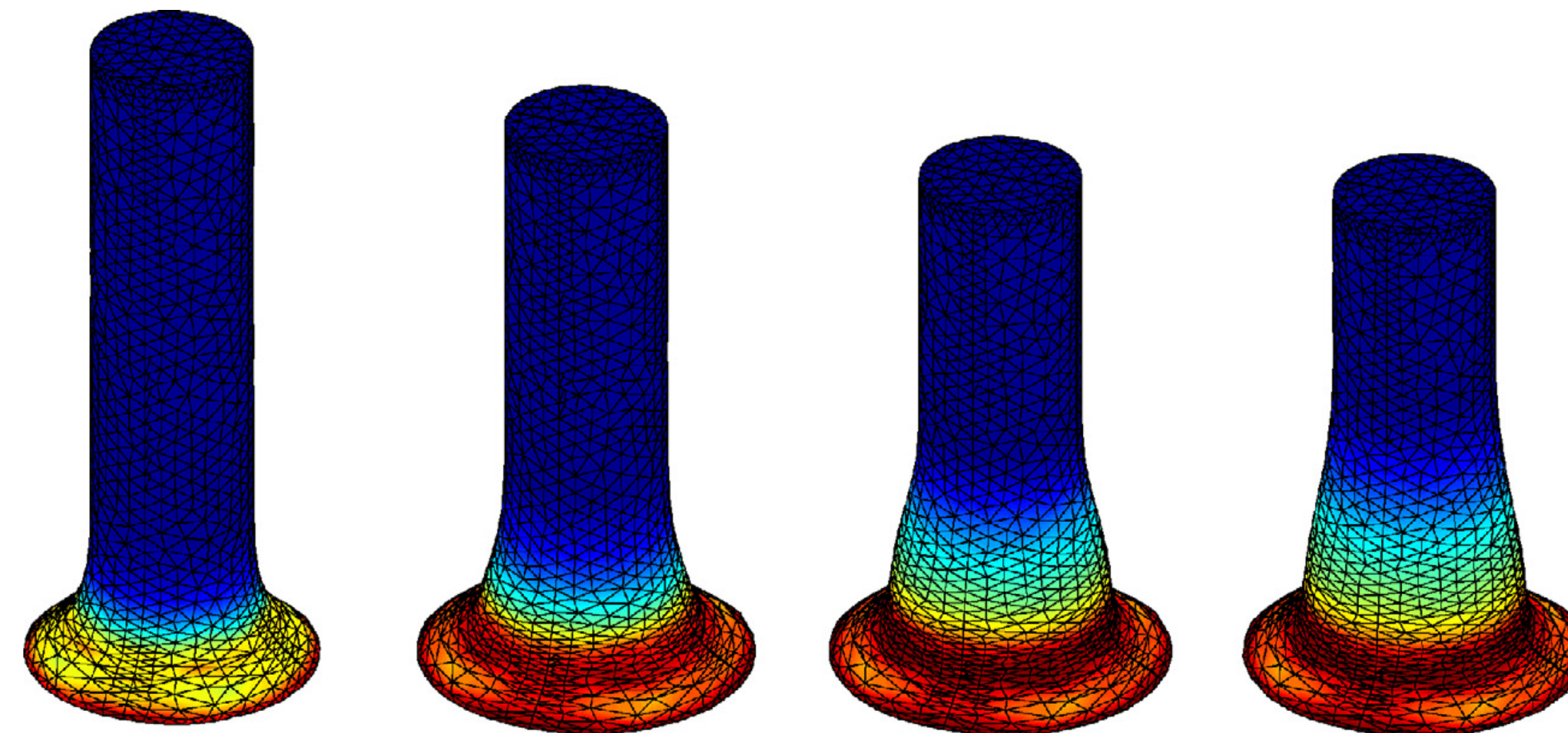


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# Impacting bar



(a) Deformed geometry predicted by the current method



(b) Deformed geometry predicted by Aguirre et al. [54]

# Summary & Conclusions

# Summary & Conclusions

OpenFOAM solvers have been developed with the **significant capabilities for solid mechanics analyses**, where strict benchmarking against industry standard commercial software has been performed.

For **linear** problems:

- **Fully implicit block-coupled procedure for linear elasticity**
- Impressive efficiency compared with standard FE approaches
- Equivalent accuracy to standard linear FE approaches

For **nonlinear** problems:

- A **large strain elasto-plastic** updated Lagrangian solver with **frictional contact** boundary conditions
- straight-forward implementation of complex nonlinear constitutive laws
- Efficient parallelisation

26<sup>th</sup> to 30<sup>th</sup> June 2016

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# References

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